

Quasi-Omnigenous Stellarator (QOS) Design Status

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Overview and Motivations

- We are developing a modest-sized experimental stellarator design based on the Quasi-Omnigenous (QO) optimization technique
- These studies are the first step towards understanding the physics of compact, low aspect ratio stellarators
- Significant progress has been made toward this goal
 - Experimental tests of these design concepts will be essential
 - Ongoing optimization effort is required to develop the concept's physics and engineering features to their full potential
- This talk will report progress in the areas of
 - coil design (preserves flux surfaces and physics properties)
 - core transport, confinement of ICRF heated populations
 - bootstrap current
 - stability

The QOS Program uses innovative methods to develop compact stellarator devices:

- Quasi-Omnigeneity (QO) is a promising approach to improving transport in stellarators
 - approximately aligns bounce-average particle drift surfaces with flux surfaces
- Can do this at relatively small aspect ratio ($A = 3-4$)
 - and maintain: good neoclassical confinement, high ballooning and kink stability limits, good energetic orbit confinement
- QOS concept exploration experiment is proposed to test the QO approach
 - will determine the optimum QO coil set for possible larger experiment
- Parameters for our scoping study are: $R_0 = 1$ m, $\langle a \rangle = 0.28$ m, $B_0 = 1-2$ T, $P_{\text{ECH}} = 0.6$ MW, $P_{\text{ICRF}} = 3$ MW, $P_{\text{LH}} = 1$ MW, and $t_{\text{pulse}} = 4$ s (at 1 T) and 1 s (at 2 T).
- Program focus: improvement of neoclassical and anomalous confinement, use of external coils rather than bootstrap current to create most of the poloidal field.

Quasi-omnigeneity is a stellarator optimization technique complementary to quasi-symmetrization methods.

- Symmetries in $|B|$ \rightarrow constant of the motion
- Quasi-symmetry \rightarrow single dominant helicity
- J^* (omnigeneity) \rightarrow longitudinal adiabatic invariant

Quasi-omnigeneity: $\langle \vec{v}_D \cdot \vec{\nabla} J^* / \omega \times \hat{n} \rangle$
 \rightarrow approximate $J^*(\cdot)$

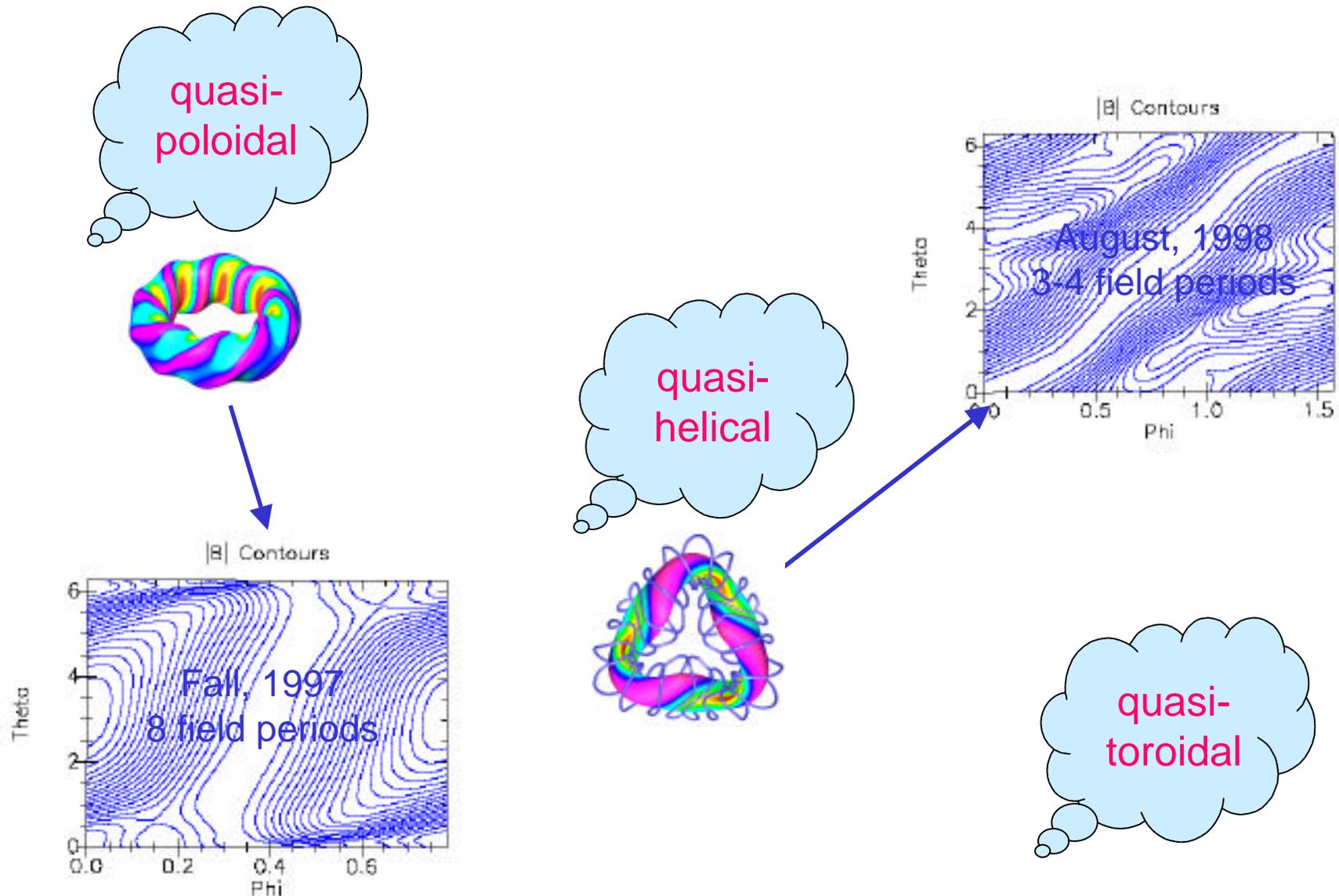
no single dominant helicity

Allowing multiple helicities provides flexibility for:

- reduction of finite bootstrap current

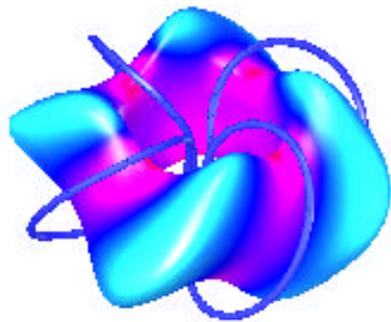
Evolution of Quasi-Omnigenous Devices

- QO configurations exist “in the neighborhood” of quasi-symmetric configurations
- Recently the QO has been used to extend nearly quasi-helical devices to low A



Stellarator optimization loop determines outer flux surface shape. Coils which produce this shape are next derived:

Initial configuration:
equilibrium from
simple modular coils



Adjust boundary
shape and
current density
profile

Solve
VMEC
equilibrium

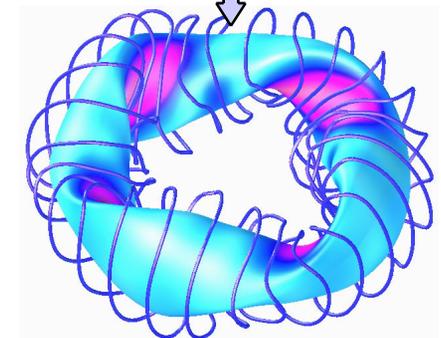
Convert to
Boozer
coords.

Calculate
2

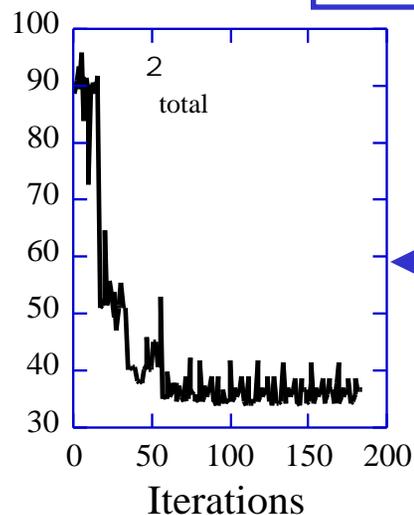
Final optimized
configuration:



NESCOIL
COILOPT

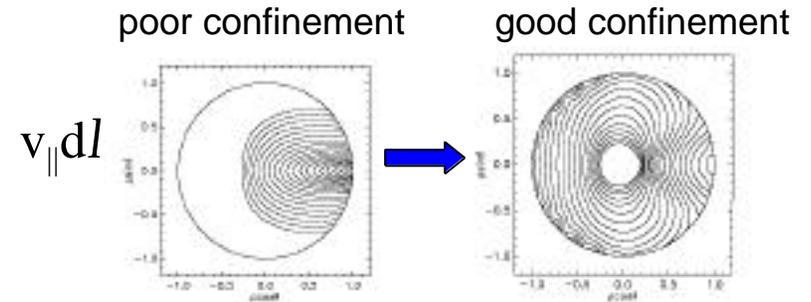


Levenberg-Marquardt
Minimize 2



QO configurations are based on a well-developed set of physics/engineering optimization targets

- Confinement is improved by aligning $J_{\parallel} dl$ surfaces (trapped orbit trajectories) with magnetic surfaces:



- Stability: maintain magnetic well, Mercier, ballooning stability
- Transform profile: $0.5 < q < 0.8$, $d q / dr > 0$ with low bootstrap current
 - avoid resonances, positive transport scaling with i , avoid neoclassical island growth
- Coil optimization (separate from outer surface optimization)
 - maintain adequate coil/plasma separation, good surface reconstruction

A New Concept Exploration Experiment, QOS, is needed to test Quasi-Omnigeneity

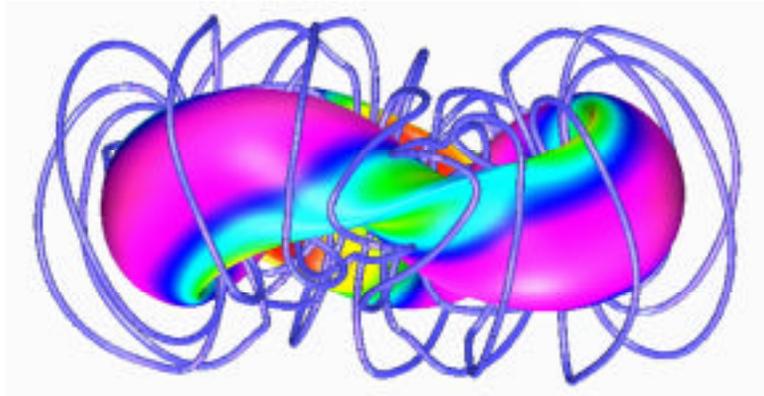
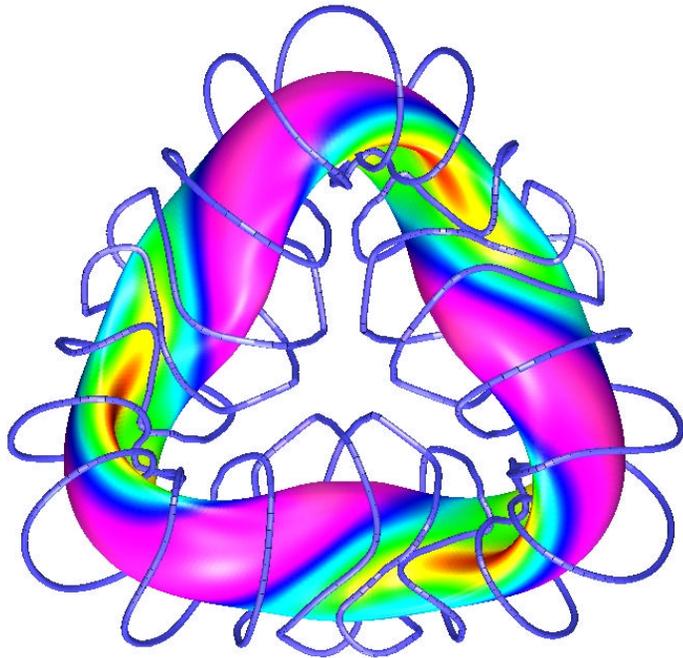
- Broaden the scientific base on the quasi-symmetry being tested in HSX and NCSX into low-aspect-ratio non-symmetric stellarators
- Test reduction of neoclassical transport via nonsymmetric quasi-omnigeneity, and the effect of radial electric fields on confinement
- Test reduction of energetic orbit losses in non-symmetric low-aspect-ratio stellarators
- Test reduction of the bootstrap current and the configuration independence on
- Test methods to affect anomalous transport, such as producing sheared $E \times B$ flow, and understand flow damping in non-symmetric configurations

Requirements for Medium-Scale Experiment to Test the QO Optimization Principle

- Adequate neoclassical confinement
 - transport reduced by factor of 2 relative to ISS95 to allow clear tests of enhanced confinement regimes
 - allow tests of improvement with ambipolar electric field
 - good reactor extrapolation
- Good confinement of energetic trapped ions
 - Viable ICRF heating efficiency
 - Adequate alpha-particle confinement for reactor extrapolation
- Bootstrap current provides < 30% of total transform
 - test configuration independence of
 - demonstrate low risk of disruptions
- Adequate size, field, and pulse length so ISS95 scaling permits tests at relevant parameters (n , T , $*$, $<$ $>$)
- Coil sets with sufficient access for heating, diagnostics
 - 3 or 4 field periods with 5-7 modular coils per period
 - $\langle a \rangle \sim 0.28$ m, $B_0 = 1 - 2$ T, pulse 1 second

Quasi-Omnigenous Stellarator

(Scoping Study Parameters)



- R_0 1.0 m
- $\langle a \rangle$ 28 cm
- $R_0/\langle a \rangle$ 3.6
- Volume 1.6 m³
- B_0 1 T (4 s)
- 2 T (1 s)?
- Plasma Heating (10 s)
 - 0.6 MW; 53.2 GHz
 - 2 MW; 6-20+ MHz
 - 1 MW; <40-80 MHz
 - 1 MW; 2.45 GHz

QOS Plasma Performance

- **Consistent Sets of Plasma Parameters***

<u>Plasma Parameter</u>	<u>0.4-MW ECH</u>	<u>1-MW ICRF</u>	<u>2.5-MW ICRF</u>
B_0 (T)	1	2	2
n_e ($10^{19}m^{-3}$)	1.6 ^(a)	3.2 ^(b)	24
τ_E (ms)	17	43	41
β (%)	0.7	0.5	2.8
T_{e0} (keV)	3.4	4.3	1.7
T_{i0} (keV)	--	--	1.7

- (a) 2nd harmonic X-mode

- (b) O-X electron Bernstein wave mode conversion

- (c) O mode fundamental

* based on ISS95 stellarator scaling with a confinement

* improvement factor $H = 2$ ($H = 1.4 - 3$ obtained in W7-AS)

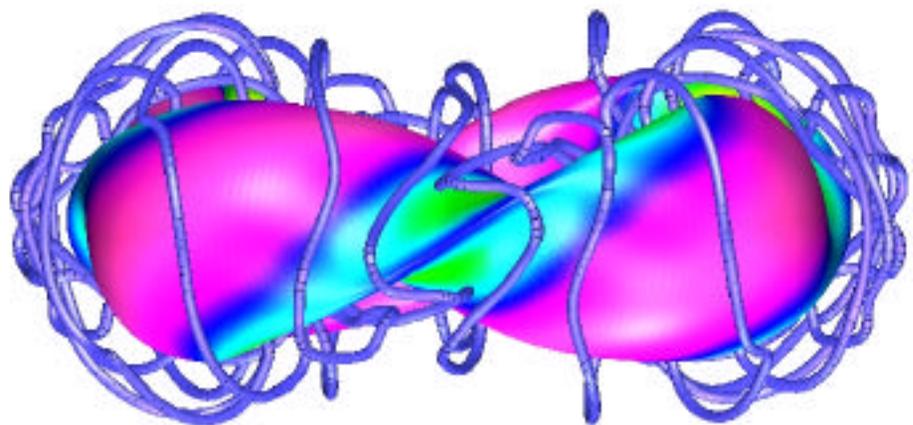
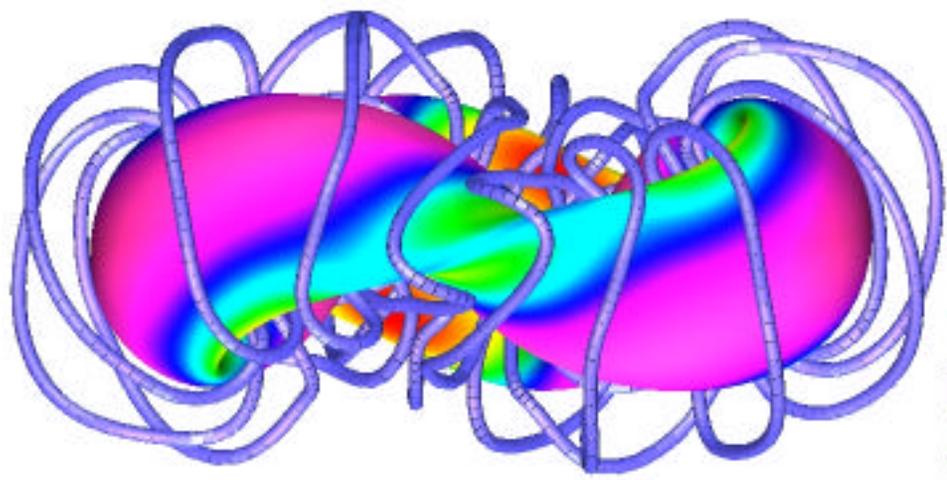
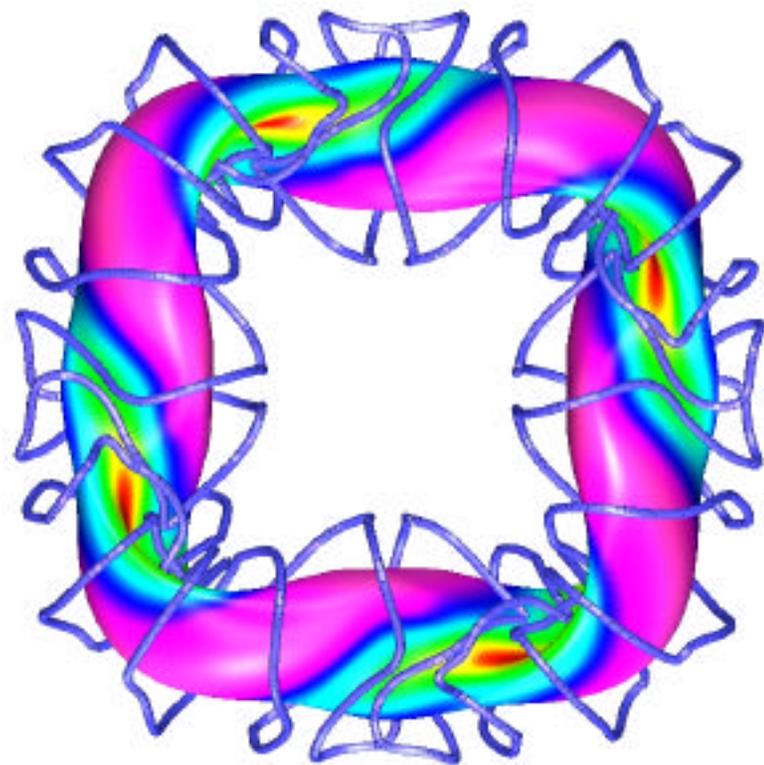
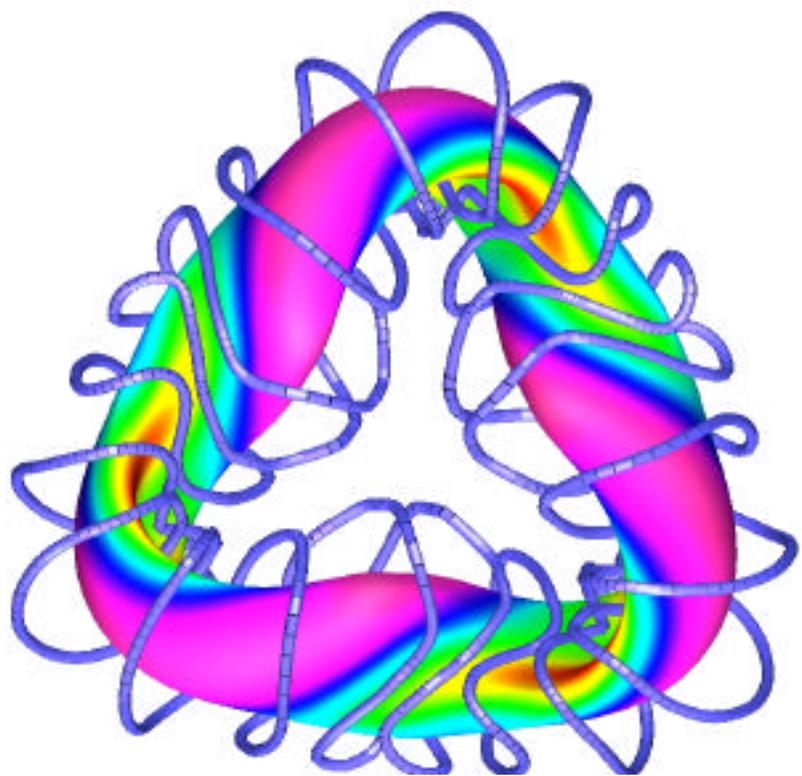
– τ_E (ISS95) = $H \times 0.079 \langle a \rangle^{2.21} R^{0.65} P^{-0.59} n^{0.51} B^{0.83} i^{0.4}$

– assumes τ_E (neoclassical) > several times τ_E (ISS95)

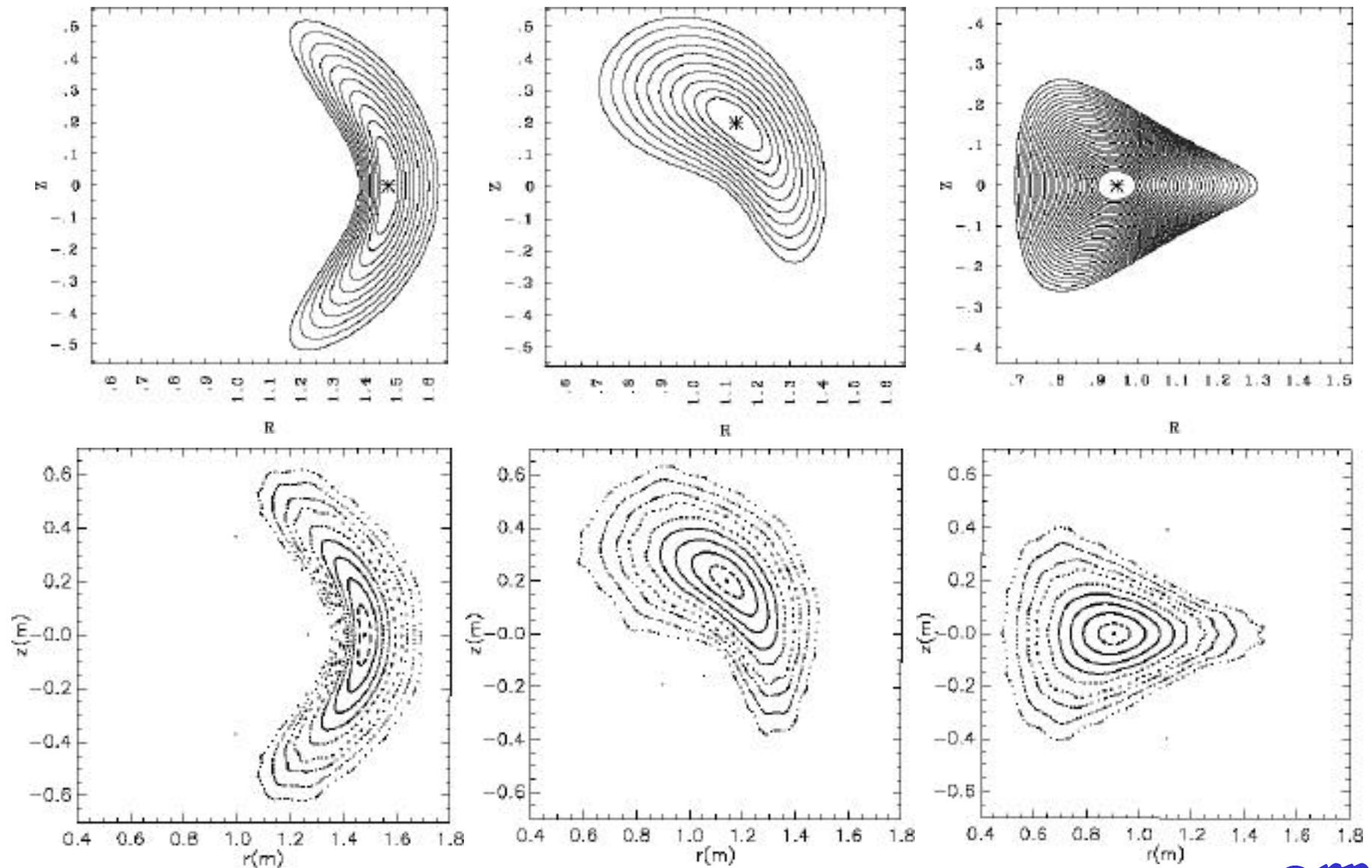
QOS Research Program Goals

- **Reduction of anomalous transport $\Rightarrow H > 2$**
 - **profile control, sheared E x B flow using ICRF**
- **Reduction of energetic orbit losses and neoclassical transport in low-A, non-symmetric configurations**
 - **modification of |B| Fourier components with auxiliary coils to spoil the optimization**
 - **modify radial electric field with ICRF, biased probes**
- **Study control of the bootstrap current**
 - **cancellation by different |B| harmonics, dependence on**
 - **test configuration independence on**
- **Study stability at low β ($< \sim 1.5\%$)**
 - **turbulence spectra, kink stability with OH driven current**

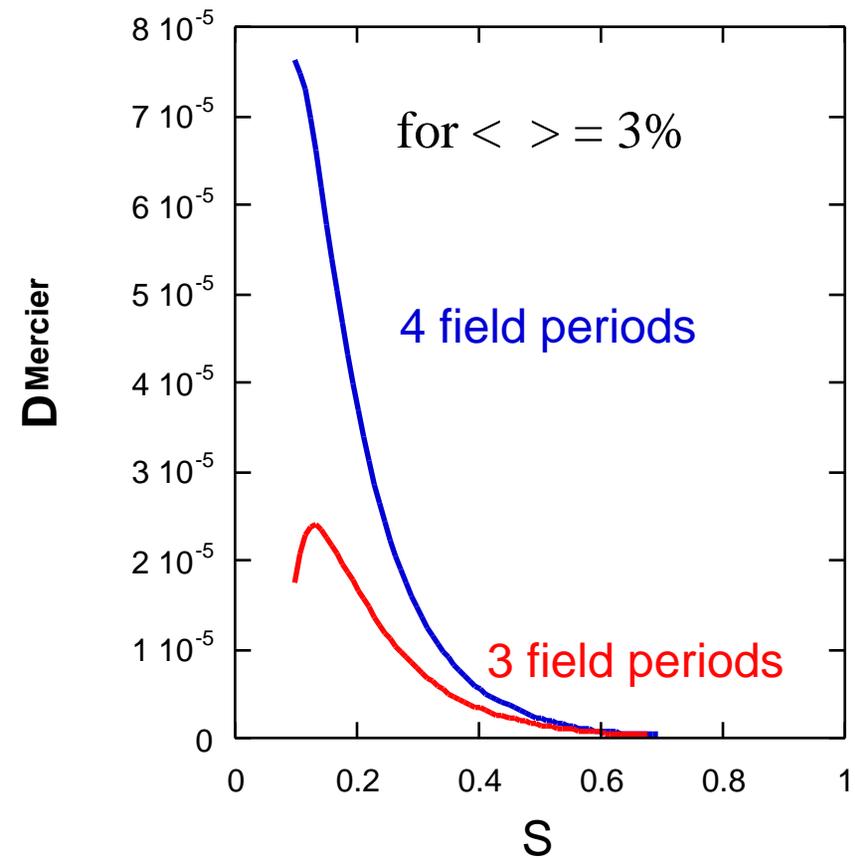
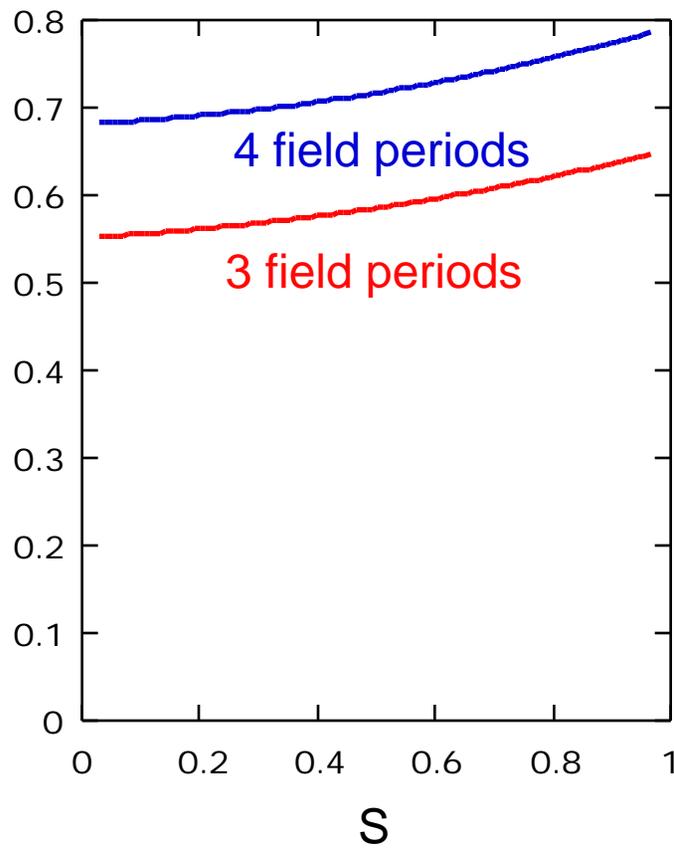
Recent 3 and 4 field period
QO configurations



Excellent flux surface reconstruction is obtained using modular coils generated by the COILOPT code:



Rotational transform and Mercier coefficient profiles for $N_{fp} = 3$ and 4 configurations

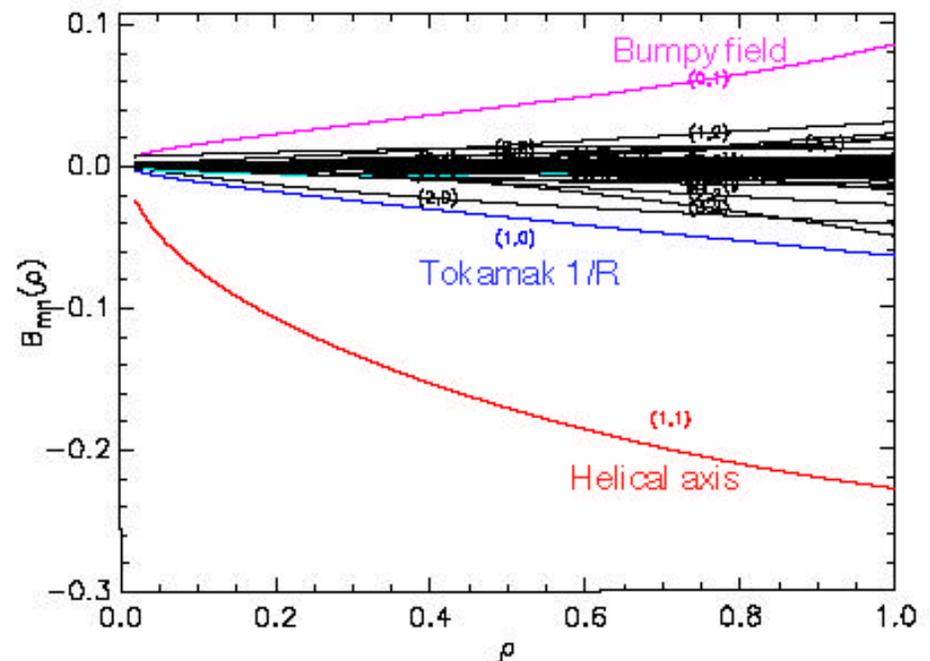
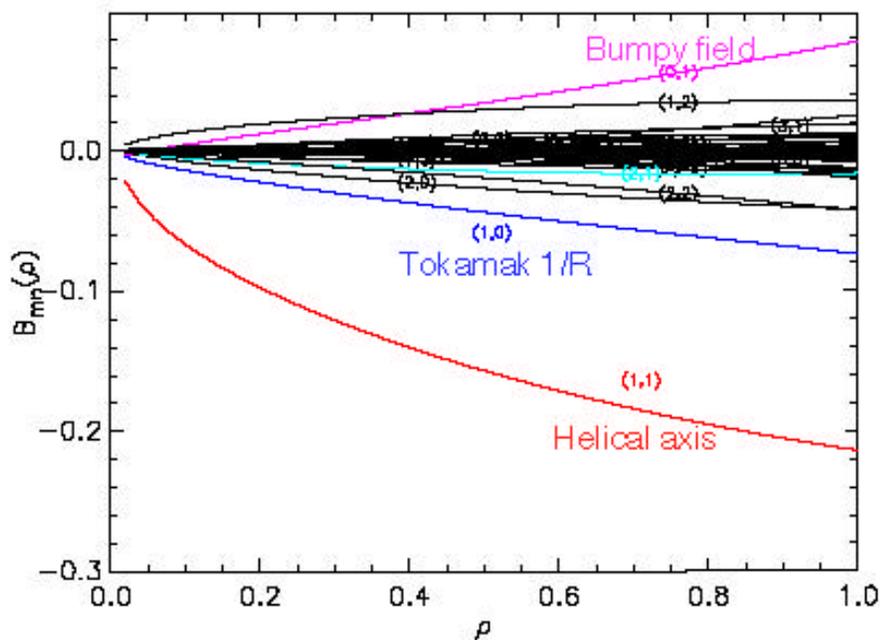


The B_{mn} spectra show that the helical component is dominant. The $1/R$ term is down from its axisymmetric tokamak level about a factor of 4.

Note: For these cases, $B_{0,0} = 1$

$$N_{fp} = 3, R_0/\langle a \rangle = 3.6$$

$$N_{fp} = 4, R_0/\langle a \rangle = 4.2$$



Effectively a hybrid of W7-X and HSX: strong helical component is like HSX, bumpy component is like W7-X

Tools used/available in the analysis of QO configurations:

- Optimization, 3D equilibrium: VMEC
- Transport, confinement
 - DELTA5D (ORNL Monte Carlo Code)
 - FAFNER-2 (IPP/CIEMAT parallel Monte Carlo code for the T3E)
 - SYMPORBIT (symplectic orbit integrator)
 - J^* , B_{\min} , B_{\max} contour plotting
- Stability
 - COBRA (fast matrix/variational 3D ballooning, R. Sanchez)
 - CHAFAR (Averaging method)
 - Resistive MHD for 3D configurations (L. Garcia)
- Bootstrap Current
 - Collisionless limit
 - NIFS multi-regime code
- Coil Design
 - COILOPT
 - NESCOIL

Energetic Collisionless Orbit Confinement

- ICRF tail populations

- ensemble of particles started out at $B = B_{\text{res}}$ locations with $v_{\parallel 0} / \mu = B_{\text{res}}$ (i.e., $v_{\parallel 0} = 0$)

- Beams

- particles born as beam ionizes on intersection with 3D flux surfaces

- Alpha -particles

- uniform distribution in $v_{\parallel 0} / v$, ,

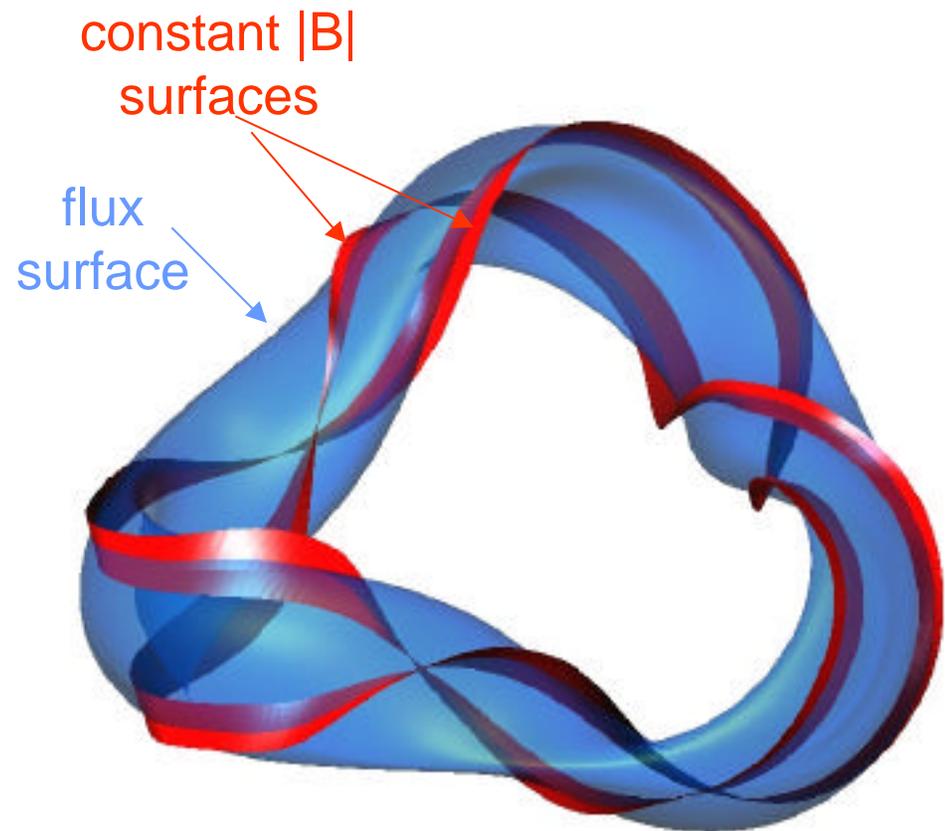
- Two issues:

- losses due to localized regions of unclosed J^* contours (all stellarators, even W7-X, Helias, etc., we have studied have this to some extent)

- deviation of energetic particle guiding center orbits away from J^* contours (becomes larger proportional to $\sqrt{\langle a \rangle}$)

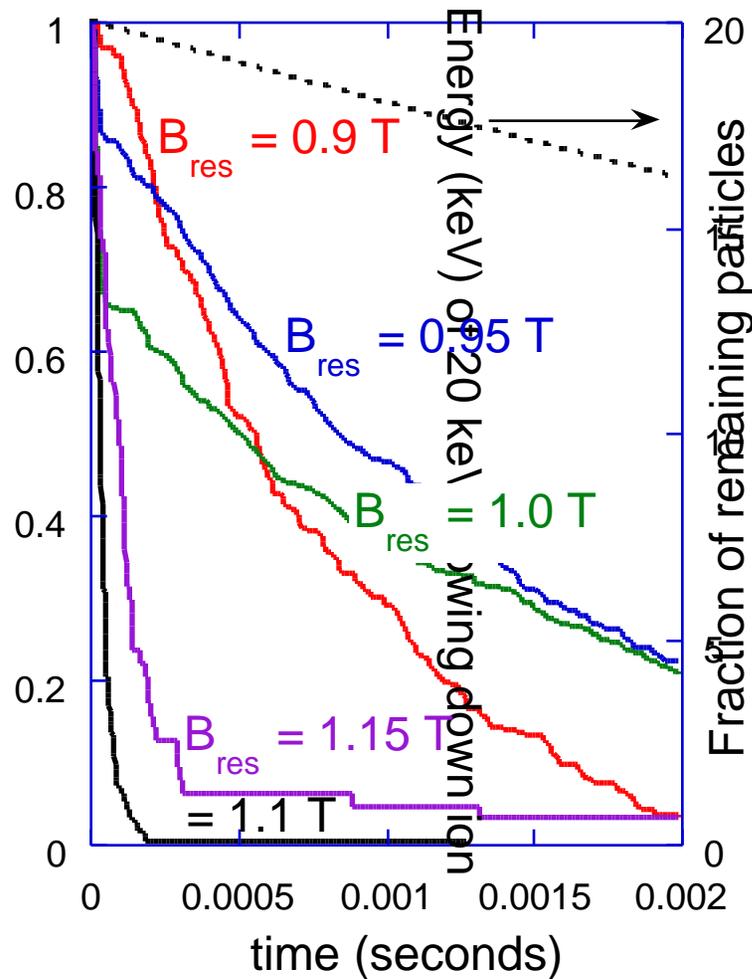
ICRF heating in stellarators depends on wave propagation, geometry of the resonant regions and the orbit confinement of the resonant ions

- Confinement of ICRF heated ions is examined by following ~500 orbits
- Intersections of $|B|$ contours with flux surfaces are determined for inner half of the plasma volume
- Ions are started out at $B = B_{\text{res}}$ with $v_{\parallel 0}/v = 0$ (equivalent to $r/\mu = B_{\text{res}}$)
- Ions leaving the outer surface are removed from the population

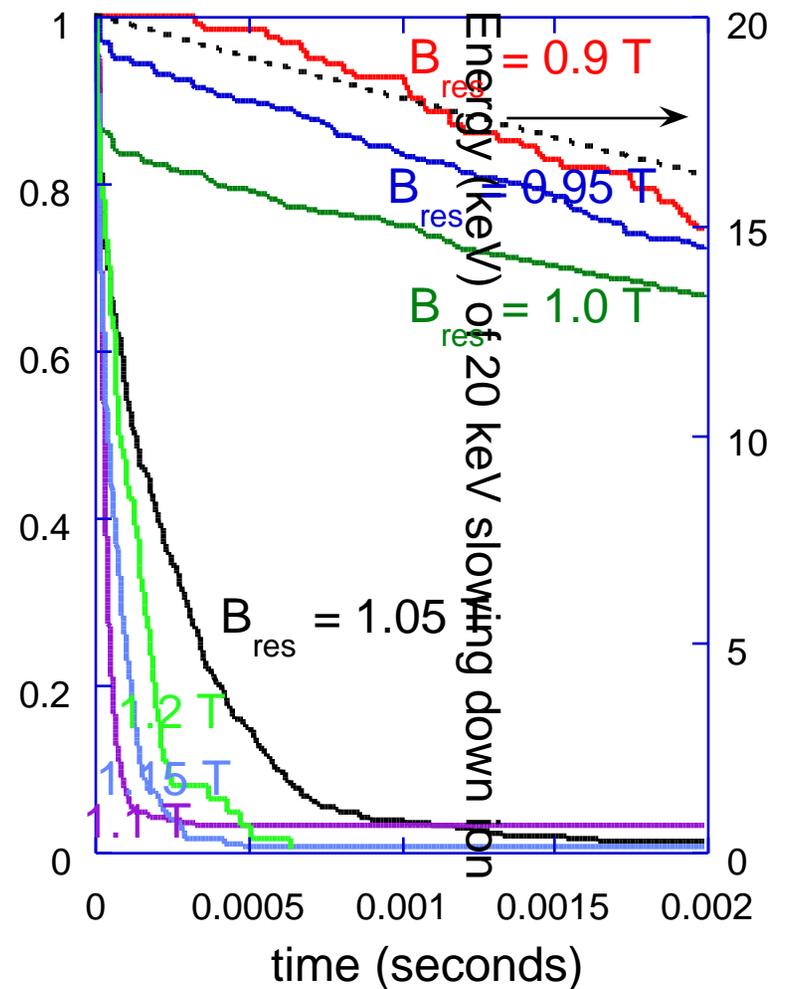


Loss rates of 20 keV ICRF ion populations (500 particles) are a sensitive function of the resonant magnetic field $B_{res} = \omega / \mu$

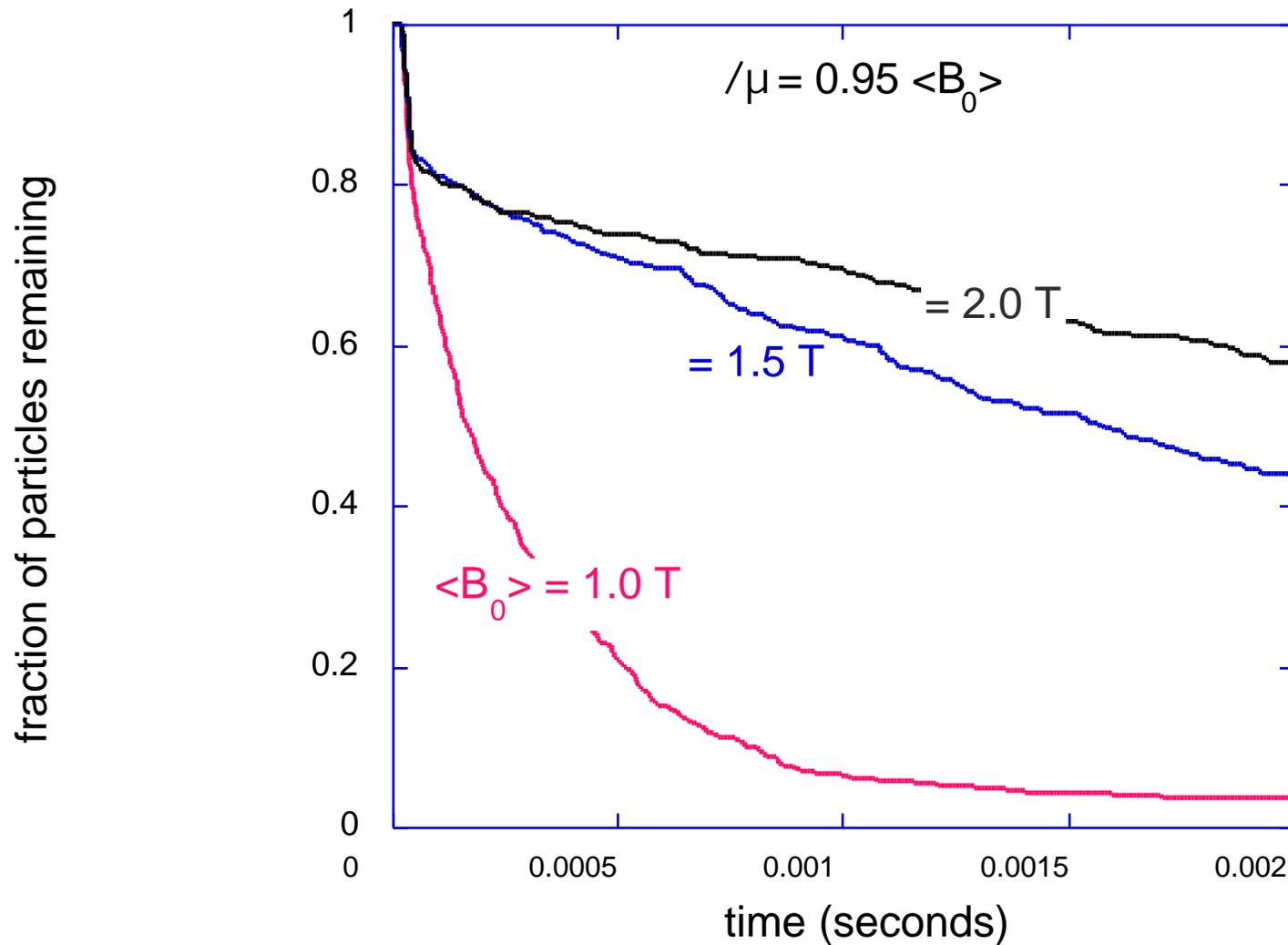
$N_{fp} = 3$



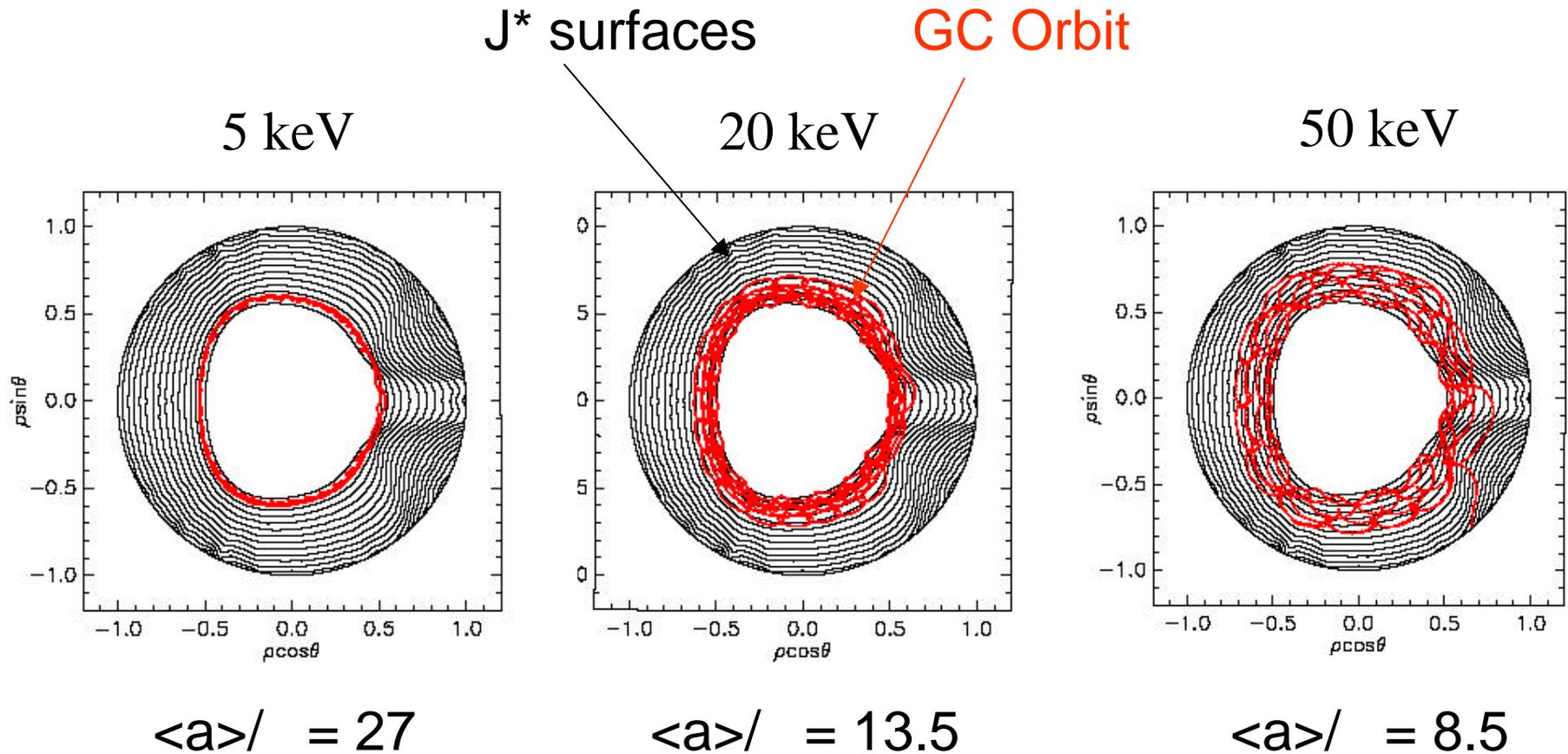
$N_{fp} = 4$



Increasing the magnetic field can significantly improve the confinement of trapped ICRF tail populations:



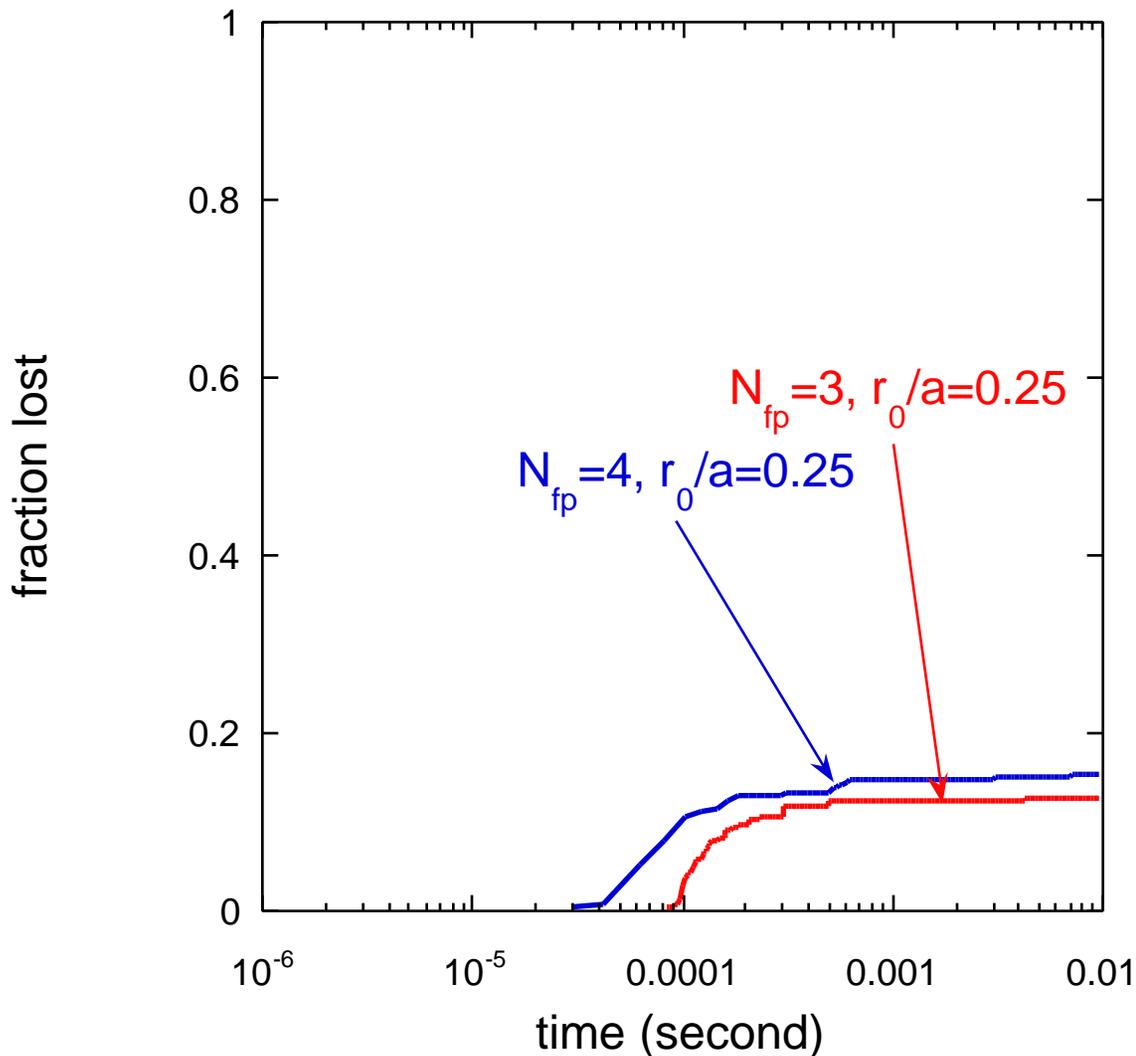
Scaling with $\langle a \rangle /$ of deeply Trapped proton orbit trajectories in a
 $B = 1\text{T } N_{fp} = 3$ QOS Device (shown in Boozer coordinates):



Confinement of heating populations in near term devices can be more demanding than alpha confinement in reactor-sized systems:

- 20 keV proton $B = 1\text{T}$ $\langle a \rangle = 24\text{ cm}$
 = 2 cm $\langle a \rangle / r = 12$
- 3.5 MeV alpha $B = 5\text{T}$ $\langle a \rangle = 2\text{ m}$
 = 5.4 cm $\langle a \rangle / r = 37$
- Prompt orbit losses are determined by:
 - Closure of J^* contours
 - degree of adiabaticity - related to size of $\langle a \rangle / r$

Collisionless α -particle losses are calculated for a reactor-scale version of the $N_{fp} = 3$ and 4 configurations (i.e., $R_0 = 10$ m, $B_0 = 5$ T, results are based on 500 α -particles per surface)



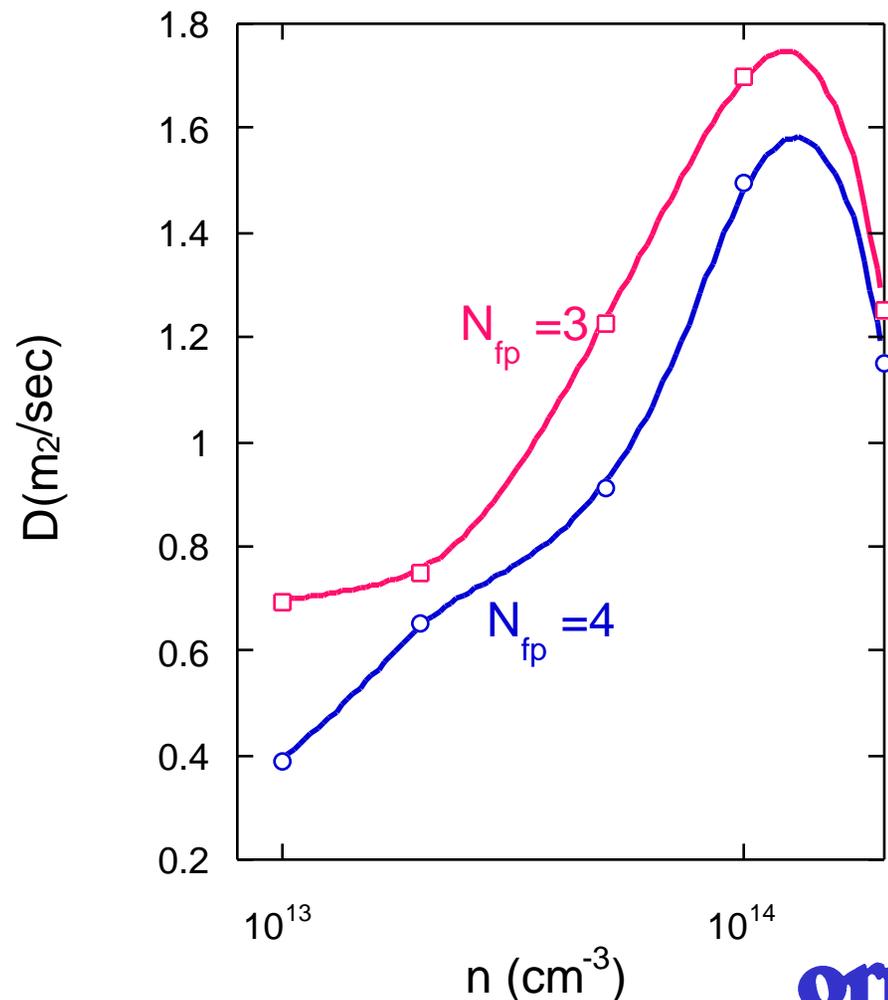
Note:

- These will be extended to $\tau_{SD} = 1$ sec with collisional effects
- This design has not yet been optimized for a reactor

Neoclassical Transport

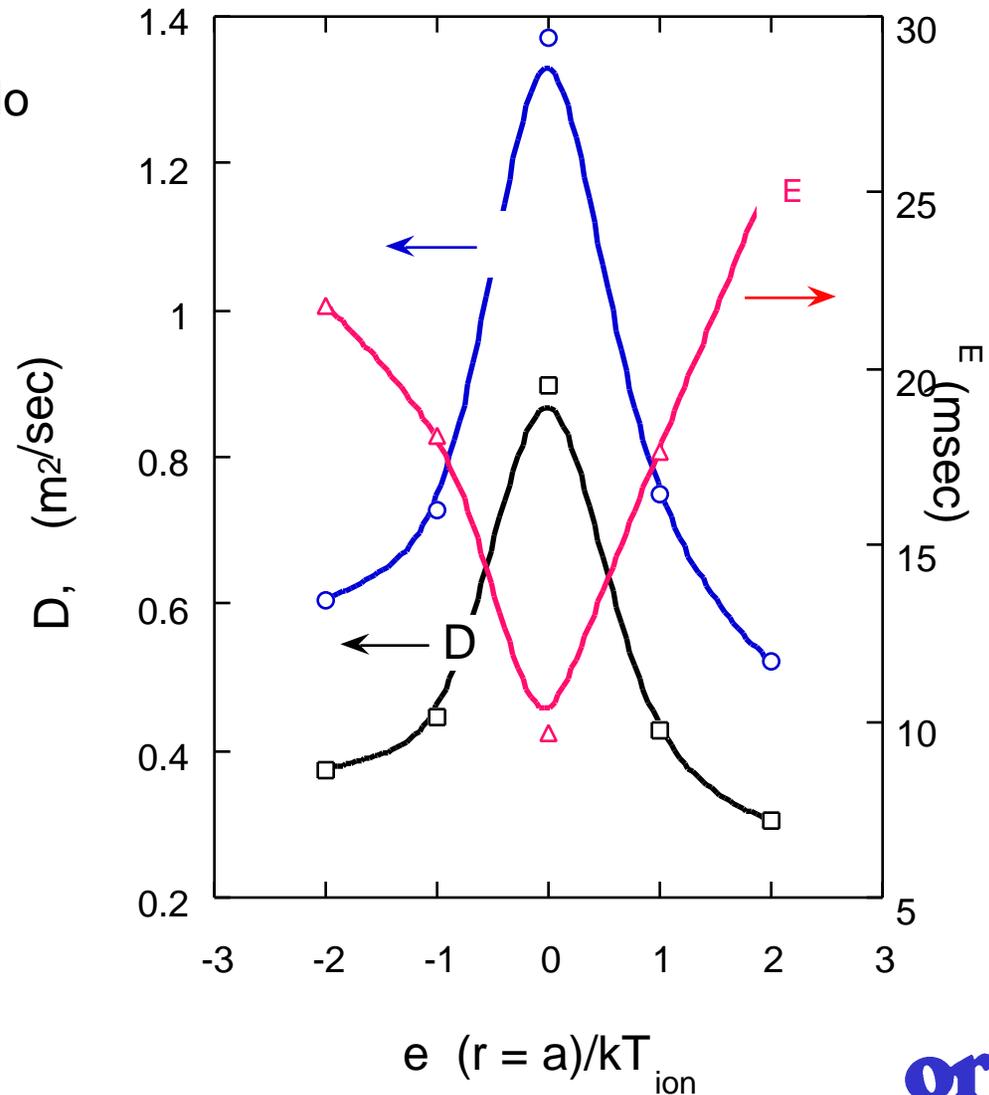
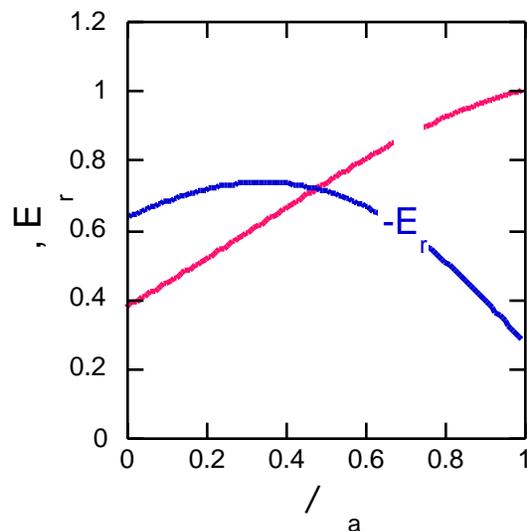
Collisionality scaling of the diffusivity for $N_{fp} = 3$ and 4 devices shows a decrease with n for $n < 10^{14} \text{ cm}^{-3}$

- $T_{\text{field}} = 1 \text{ keV}$, $E_{\text{test}} = 2 \text{ keV}$,
 $Z_{\text{eff}} = 1$, (monoenergetic, 2000 particles)
- No ambipolar electric field
- $N_{fp} = 3$
 - $0.006 < \frac{\text{eff}}{b} < 0.115$
 - $34 < L^* < 680$
- $N_{fp} = 4$
 - $0.004 < \frac{\text{eff}}{b} < 0.08$
 - $42 < L^* < 835$
- where $L^* = \lambda / L_c$, $L_c = R_0 / v$,
 $\lambda =$ mean free path, $\frac{\text{eff}}{b} = \frac{\nu_{\text{eff}}}{\omega_b}$,
 $\omega_b =$ bounce frequency



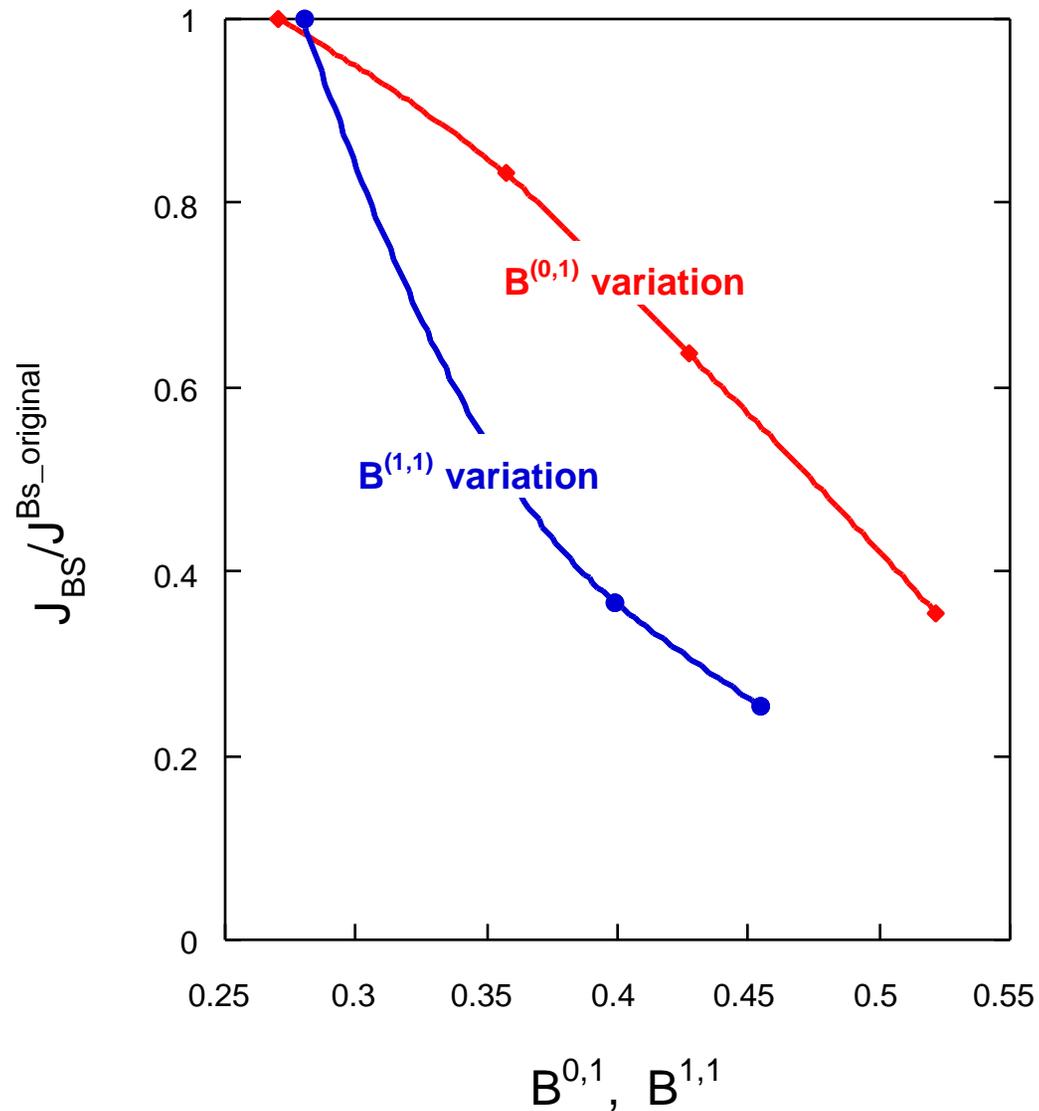
The diffusivities and 0-D energy confinement time show a strong dependence on electric field with $E \approx (2-3) \times I_{SS95}$ for typical ambipolar potentials

- $N_{fp} = 4$ configuration
- 4 monoenergetic ion Monte Carlo groups used (2000 particles each) to construct diffusivities for a Maxwellian
- $T_{field} = 1$ keV, $E_{test} = 0.5, 1, 2, 3$ keV, $n = 5 \times 10^{13} \text{ cm}^{-3}$, $Z_{eff} = 1$
- Starting position: $r/a = 0.25$ ($r/a = 0.5$)

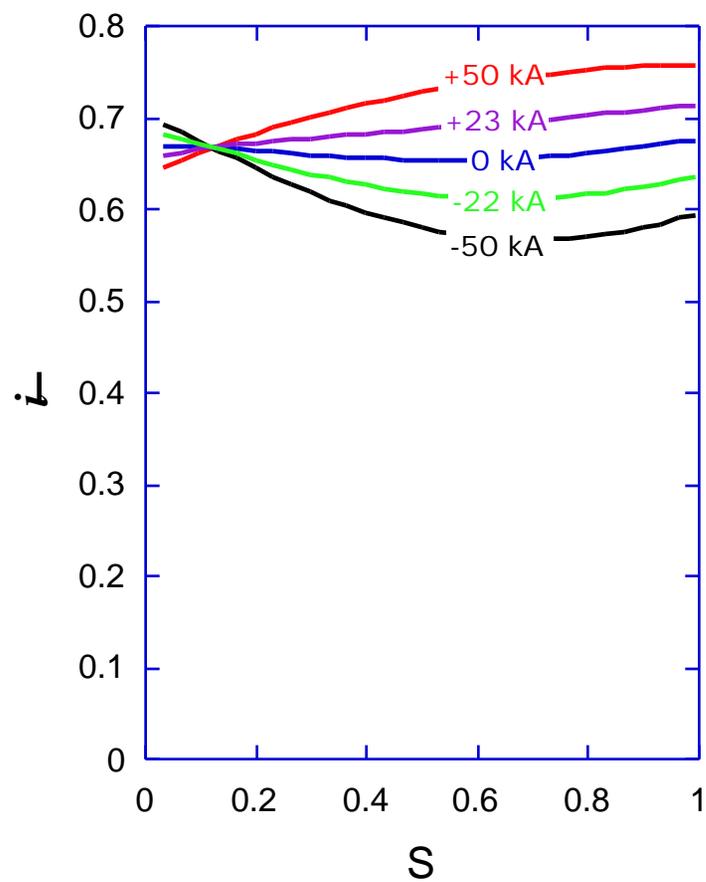
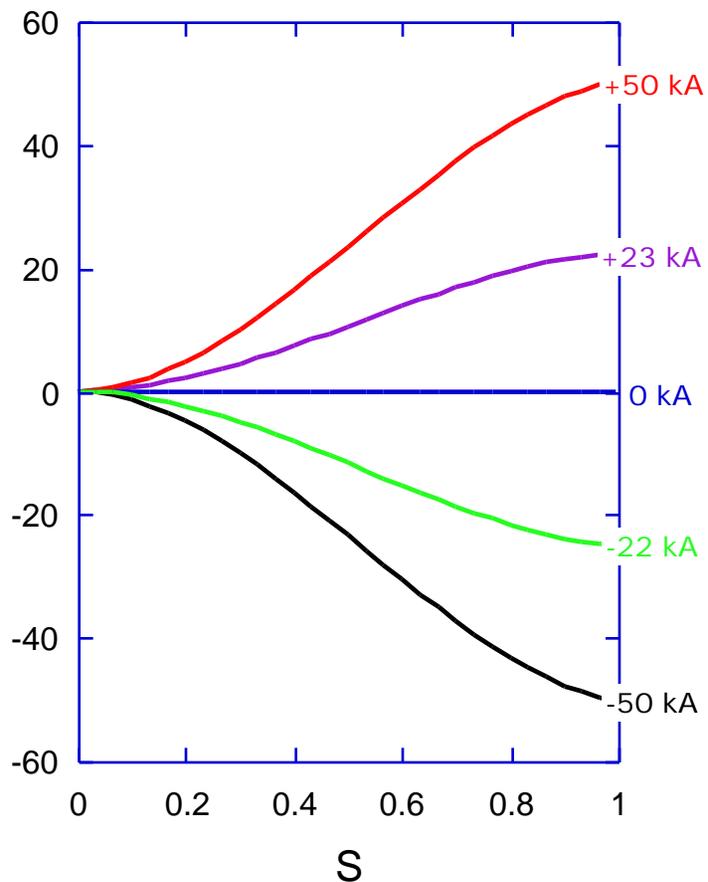


Bootstrap Current

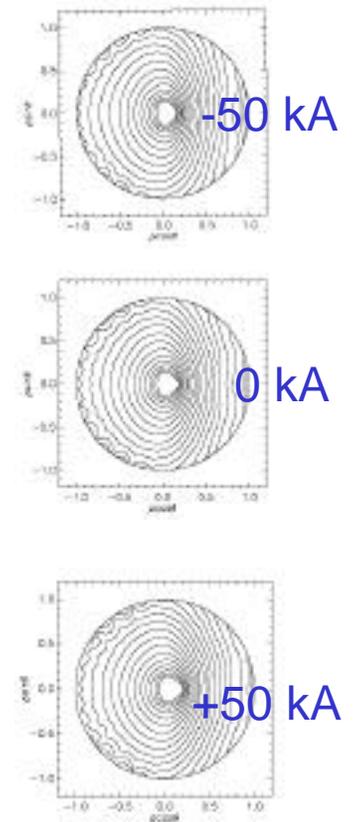
QO configurations allow flexibility with respect to the level of bootstrap current



The transform and J^* surfaces of the 3 field period QO configuration are not strongly modified by various levels/signs of bootstrap current



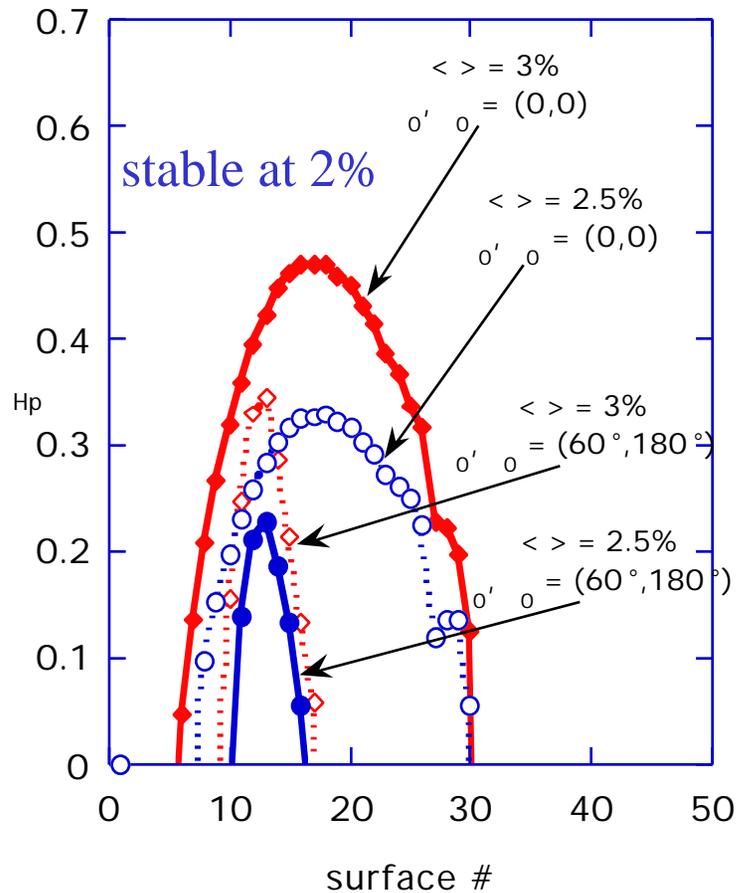
J^* contours
at $/\mu = 1$



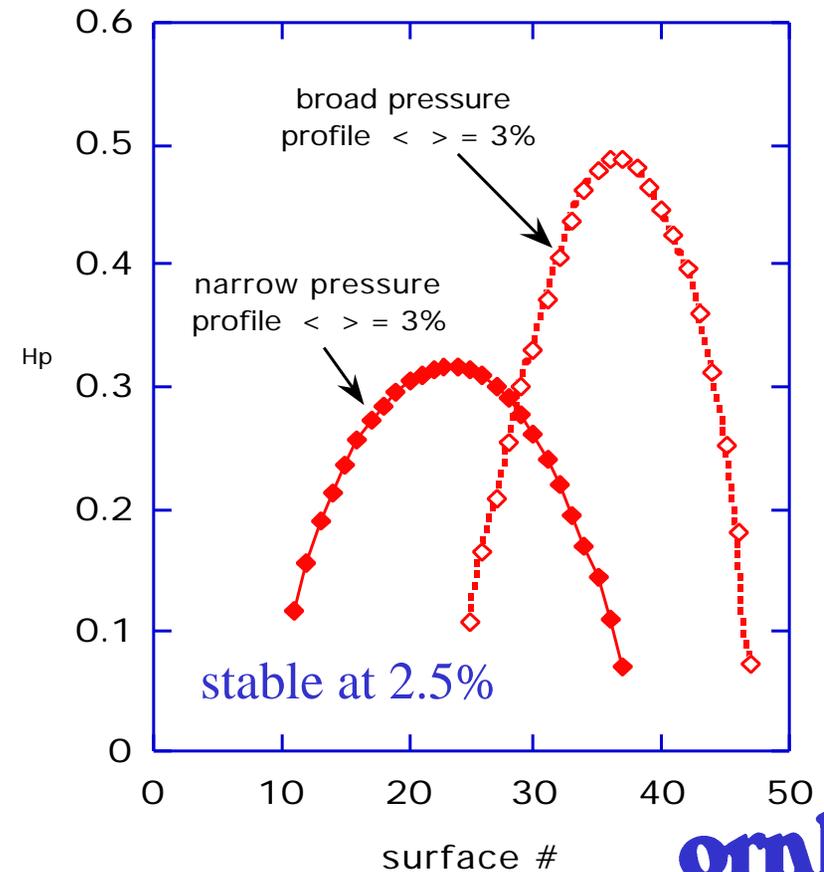
Stability

Ballooning growth rates can be a sensitive function of the pressure profile and matching point position along the field line. Recent pressure profile optimizations for the 3 period device have gotten the limit up to 3%

$$N_{fp} = 3$$

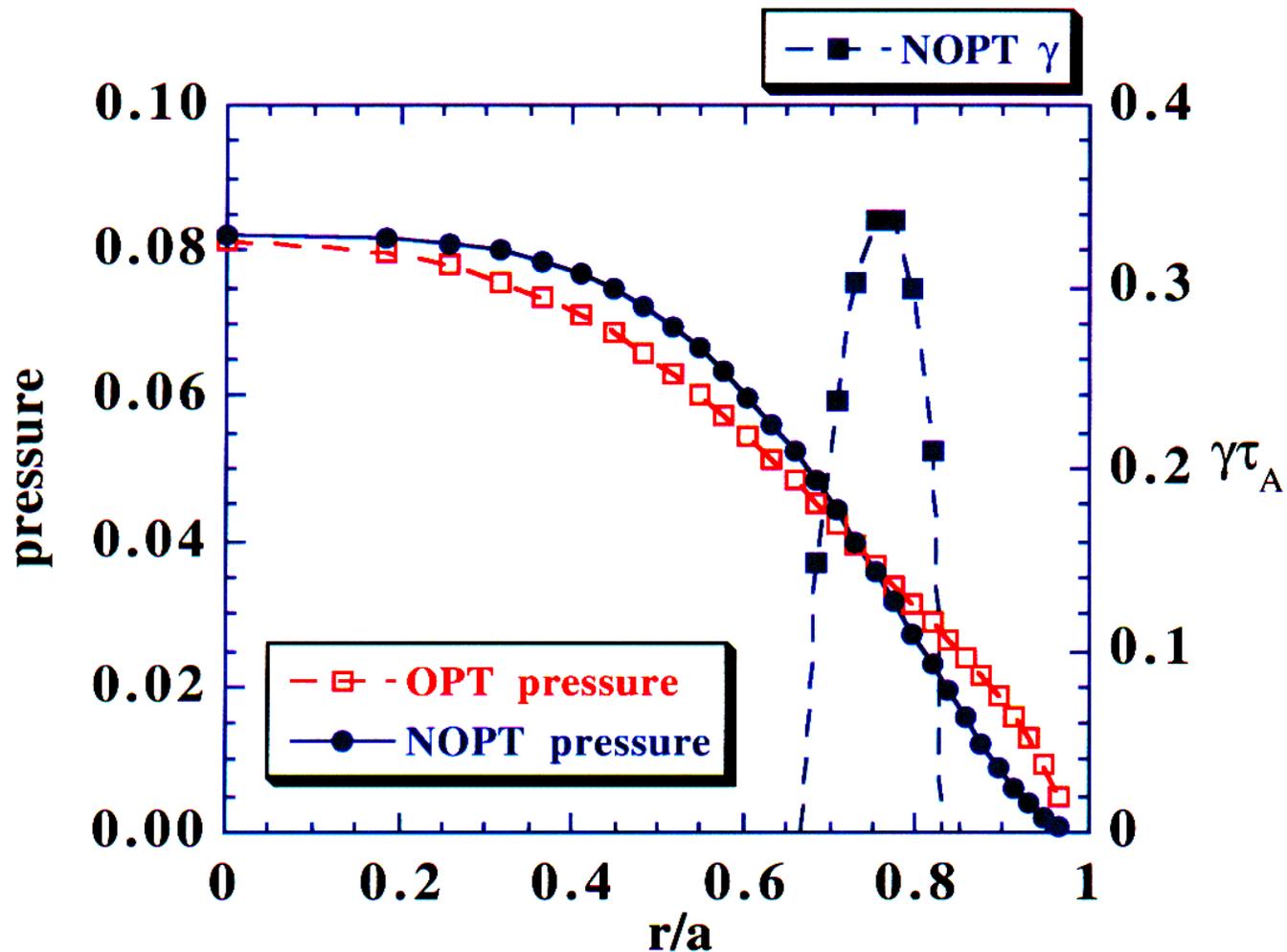


$$N_{fp} = 4$$



The first step of ballooning optimization (profile modification) has recently been completed, leading to ~20% increase in the stable :

QOS CONFIGURATION: $M=3$, $\langle\beta\rangle=3\%$, $A=3.2$



Areas for Further QOS Development

- **Optimization**: self-consistent bootstrap current and ballooning stability improvement; evaluate tradeoffs of performance vs. cost
- **Neoclassical transport**: further improvement to allow better exploration of improved confinement regimes?
- **Fast ion confinement**: sufficient for bulk ICRF heating
 - better confinement of energetic perpendicular ions needed to explore minority species heating?
- **Beta limit**: increase $<$ $>$ limit above 3% for better experimental test of reduction of bootstrap current and of beta limit dependence on magnetic configuration properties?
- **Modular coils**:
 - 5 instead of 7 coils per period for better experimental access?
 - reduce minimum bend radius for easier fabrication?

CONCLUSIONS

- Progress has been made with the preliminary design phase for a small scale QO concept exploration experiment
 - 5-7 modular coils per period - good flux surface reconstruction
 - good neoclassical transport ($E_{,neo} \approx 2-3 \times E_{,ISS95}$)
 - confined ICRF tail - efficient heating
 - can suppress/control bootstrap currents
 - ballooning $< > 2 - 3\%$ - appropriate range for experimental tests
- Unique features of QO devices
 - good test-bed for RF heating
 - optimized for confinement of high v trapped particles
 - high field access
 - can be designed to access $0.5 < i < 1$
 - rotational transform comes predominantly from coils