

The Columbia Non-neutral Torus: Status, Plans, and Opportunities for Theory Research

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With

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Columbia University
In the City of New York

CNT



CNT is a simple and compact stellarator



Magnetic surfaces are created from only four circular coils^{1,2}

$\langle R \rangle = 0.28$ m, $\langle a \rangle = 0.15$ m

Ultrahigh vacuum

$B = 0.3$ T for 15 s

Steady state at $B \leq 0.06$ T

¹ Gourdon et al., Plas. Phys. Contrl. Nucl. Fus. Research p. 849 (1969)

² Pedersen et al., Fusion Sci. Tech. 46 p 200 (2004)

Some of the questions we address in CNT

- Does a pure electron plasma have a stable equilibrium in a stellarator? (Theory: Yes^{a,b}) What is the equilibrium like? ^c
- What are the principal transport mechanisms?
- Can confinement be excellent?
 - Neoclassical transport predicts excellent confinement in a non-neutral stellarator:
 - $\tau \sim 1/\nu (e\phi/T)^2$
 - Small Debye length pure electron plasma has $(e\phi/T)^2 \approx (a/\lambda_D)^4 \gg 1$
 - So extremely long confinement is predicted, minutes or hours
- What are the properties of partially neutralized plasmas? (These can be created in a stellarator)
- Can electron-positron plasmas be created in a non-neutral stellarator^d?

a Pedersen and Boozer, PRL 88, 205002 (2002)

b Boozer, Phys. Plasmas 11, p. 4709 (2004)

c Lefrancois et al, Phys. Plasmas 12, p. (2005)

d Pedersen et al., J. Phys. B 36, p. 1039 (2003)

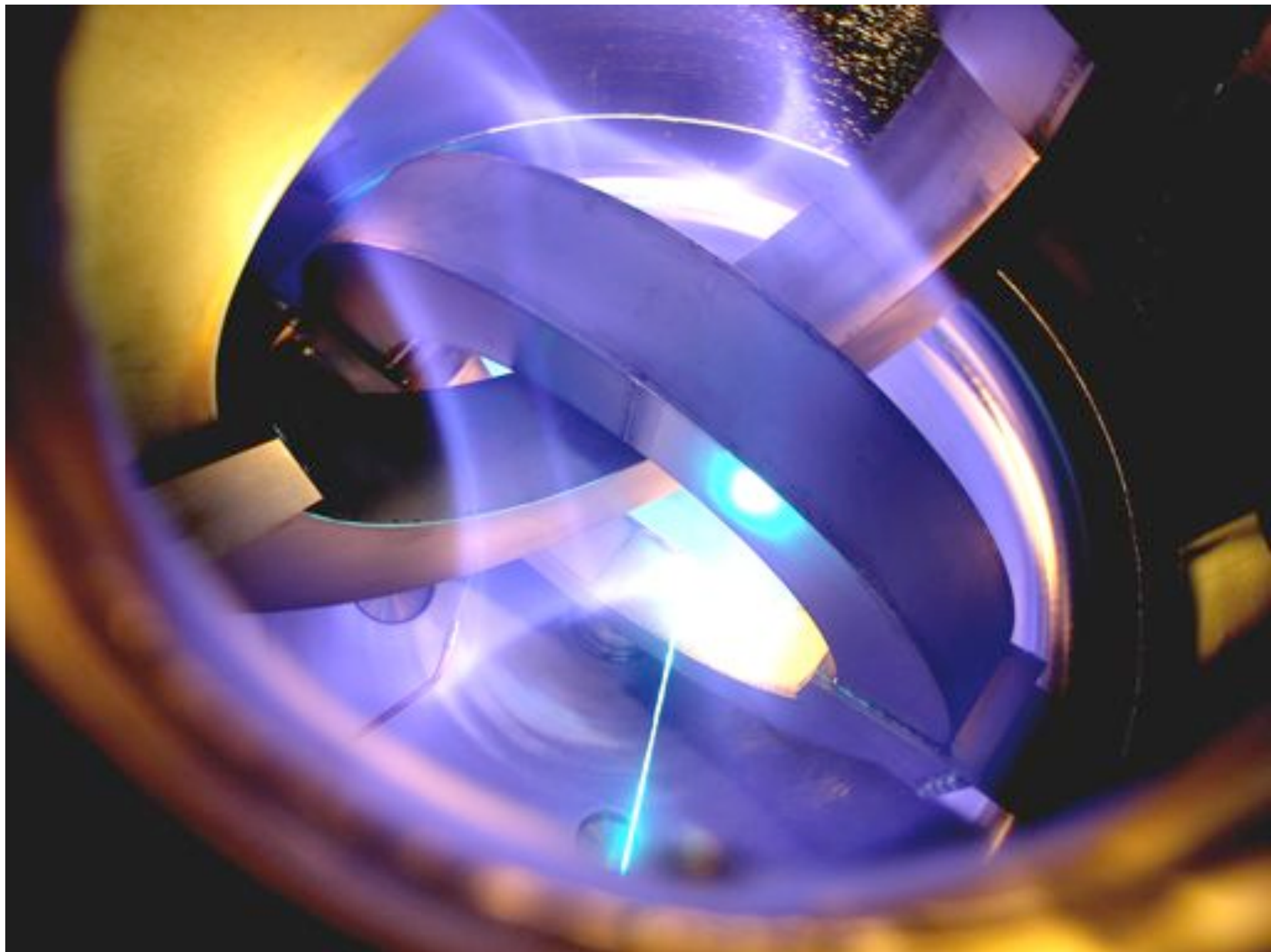
CNT fully operational since Nov 12 2004



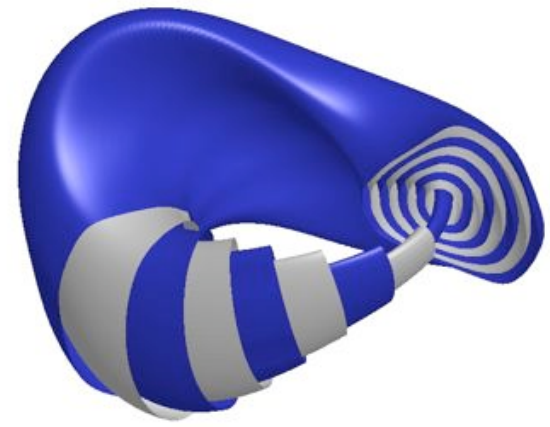
CNT's magnetic topology



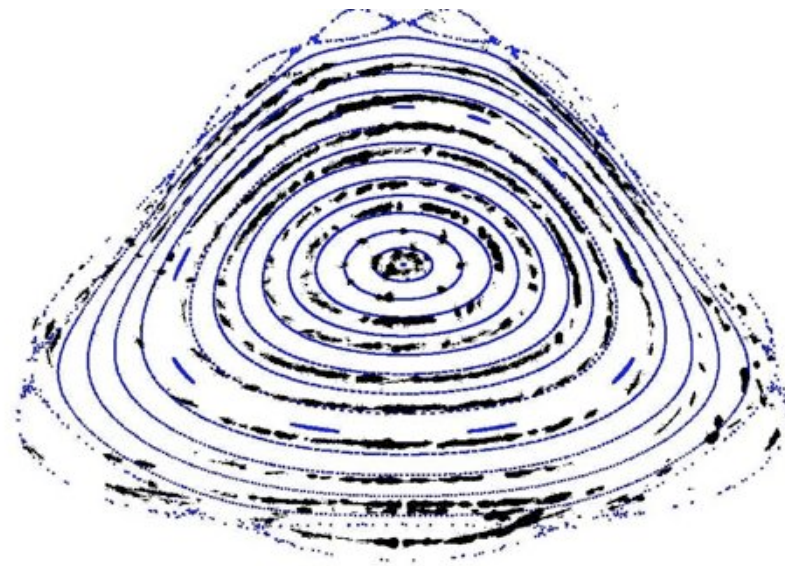
Two-period, ultralow aspect
ratio ($A < 1.9$) stellarator



Field line mapping



Field line mapping - quantitative measurements

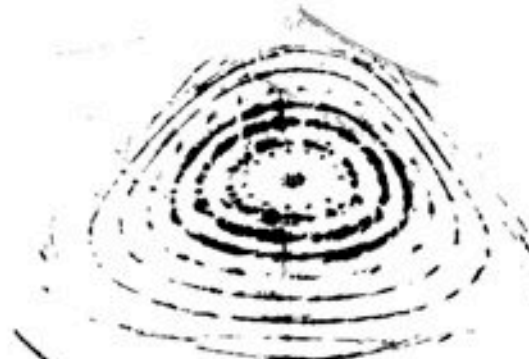


Excellent quality nested magnetic surfaces even at low B

0.10 T



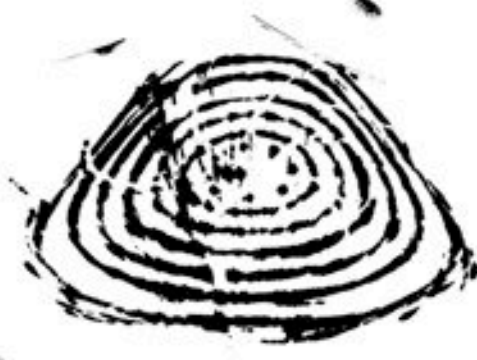
0.060 T



0.030 T



0.015 T



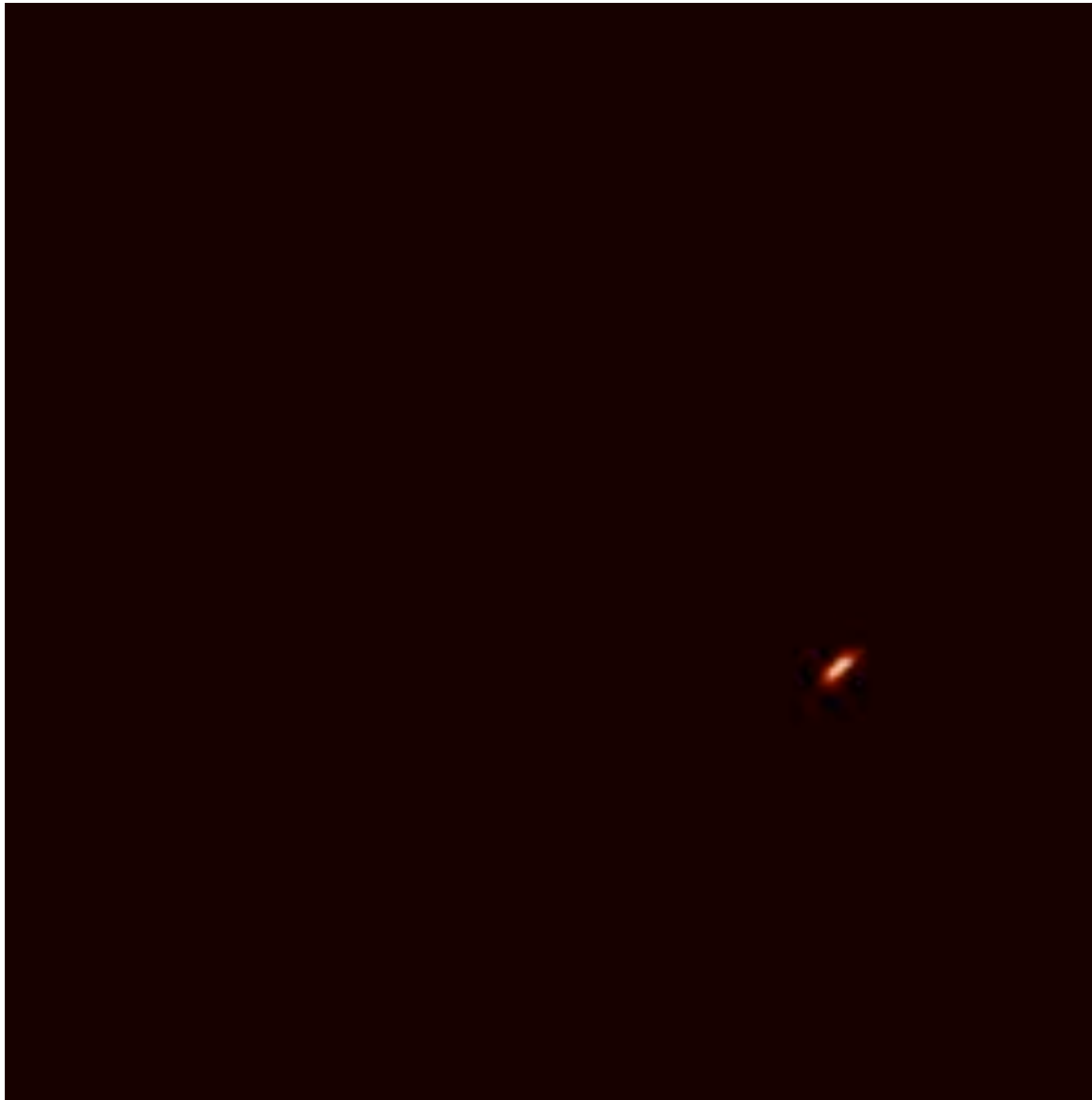
0.0074 T



0.0041 T



Present scheme for creation of pure electron plasmas



Electron source:

Thermionic emission from
heated tungsten filament

Placed on magnetic axis

Biased negatively

No anode (unlike e-gun)

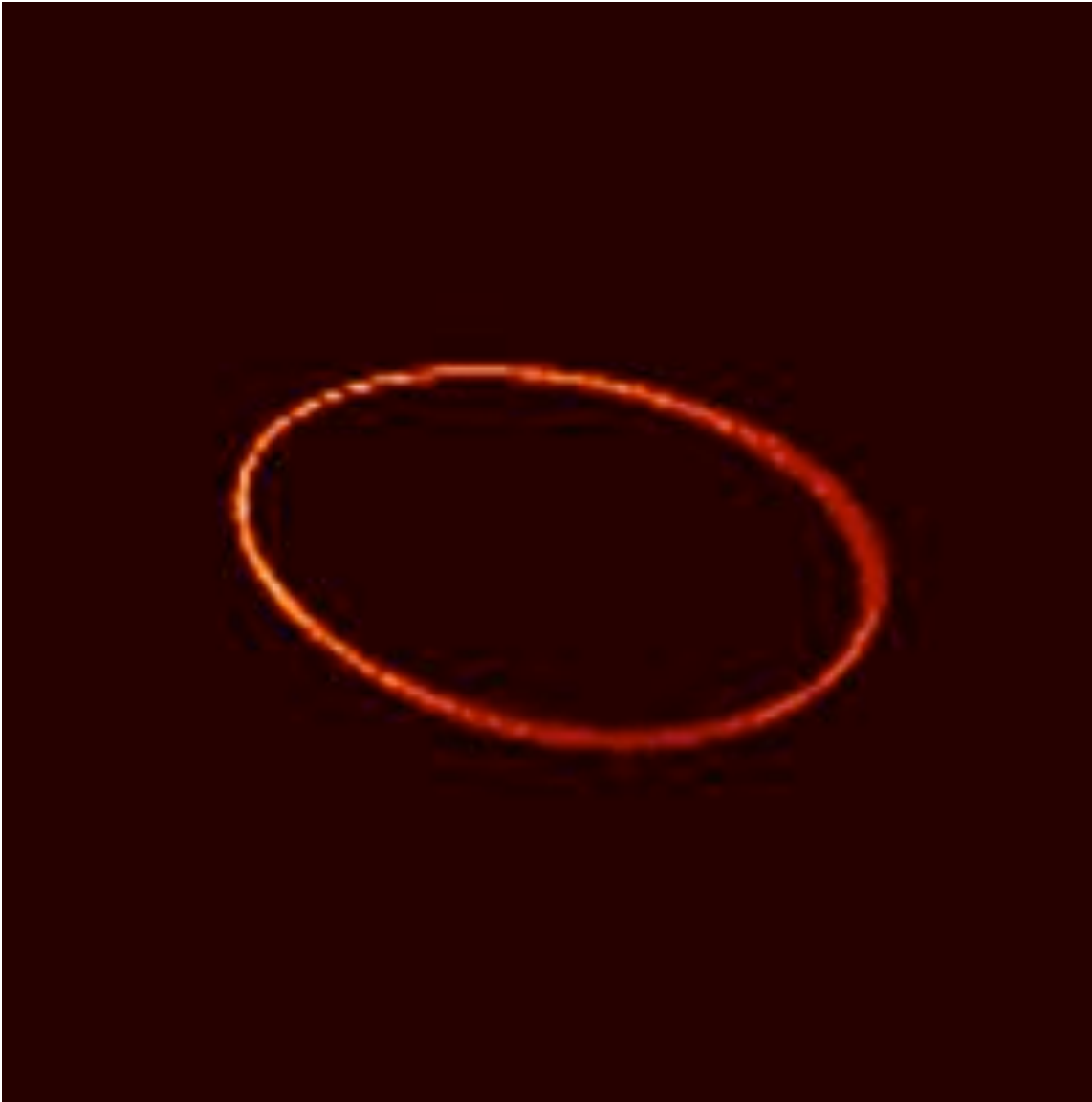
Parallel transport fills field
line on axis in $\sim 1 \mu\text{s}$

Perpendicular transport
fills the rest of the
surfaces

Reach steady state
between emission and
radial losses



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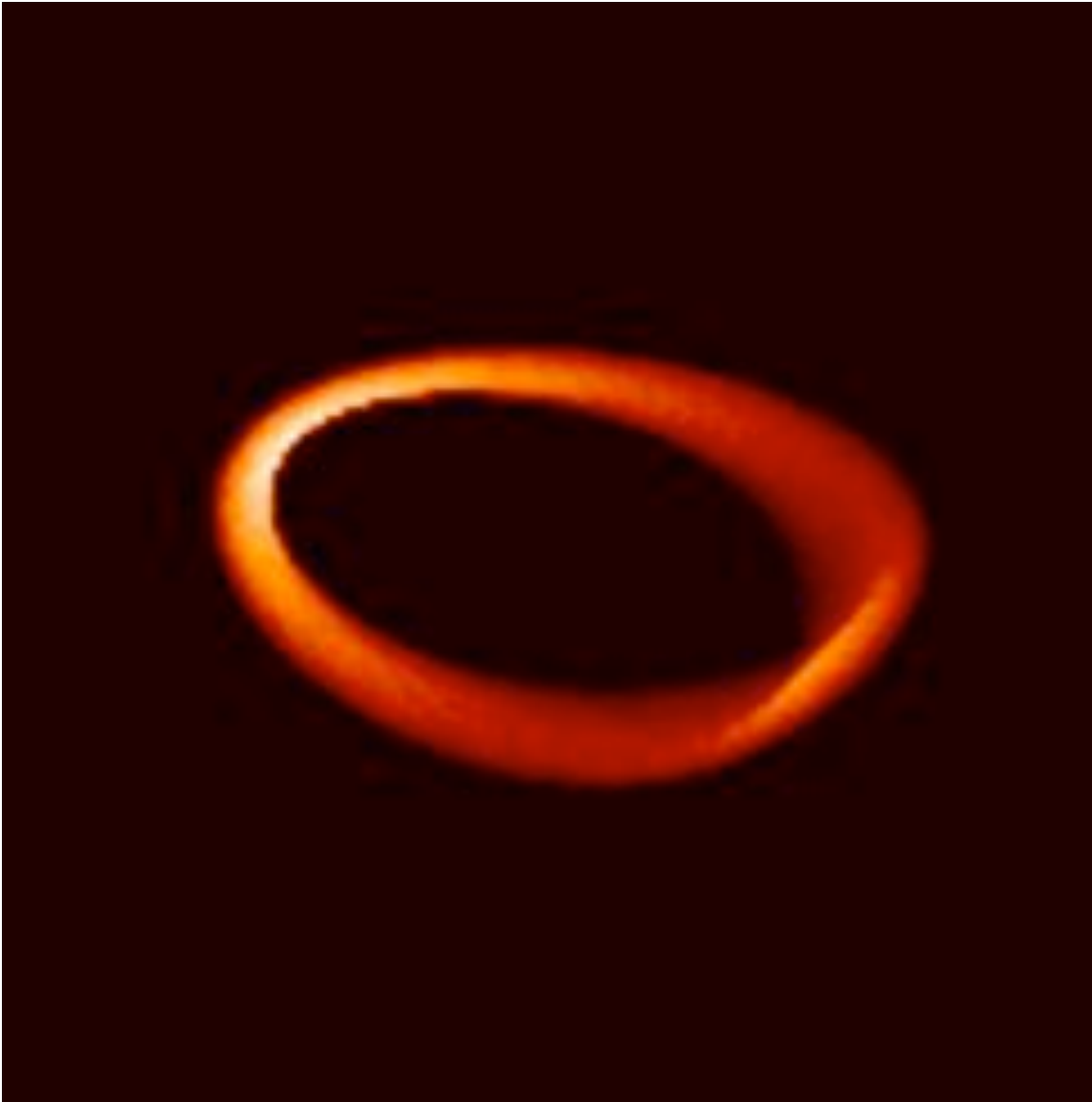
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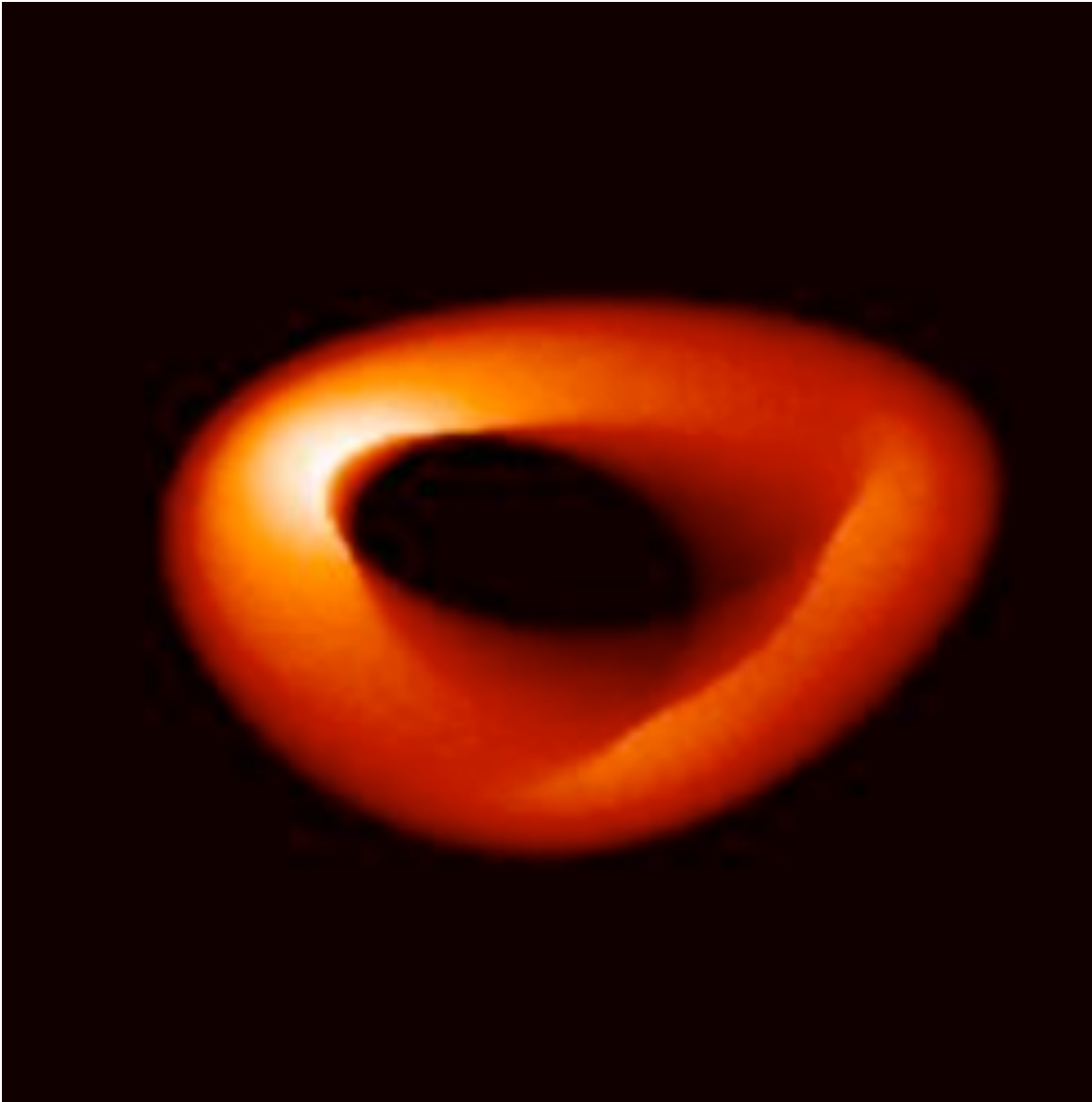
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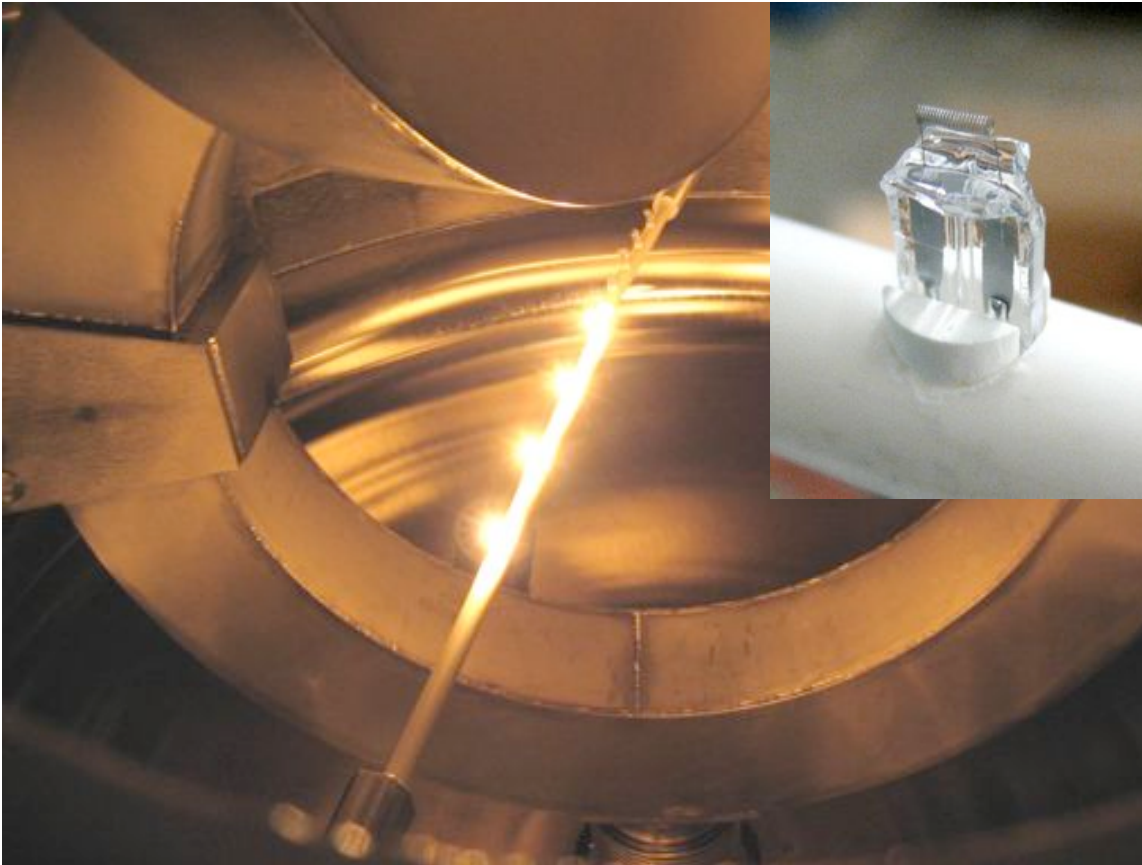
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Emitting filaments serve multiple functions



2 ceramic rods with arrays of filaments

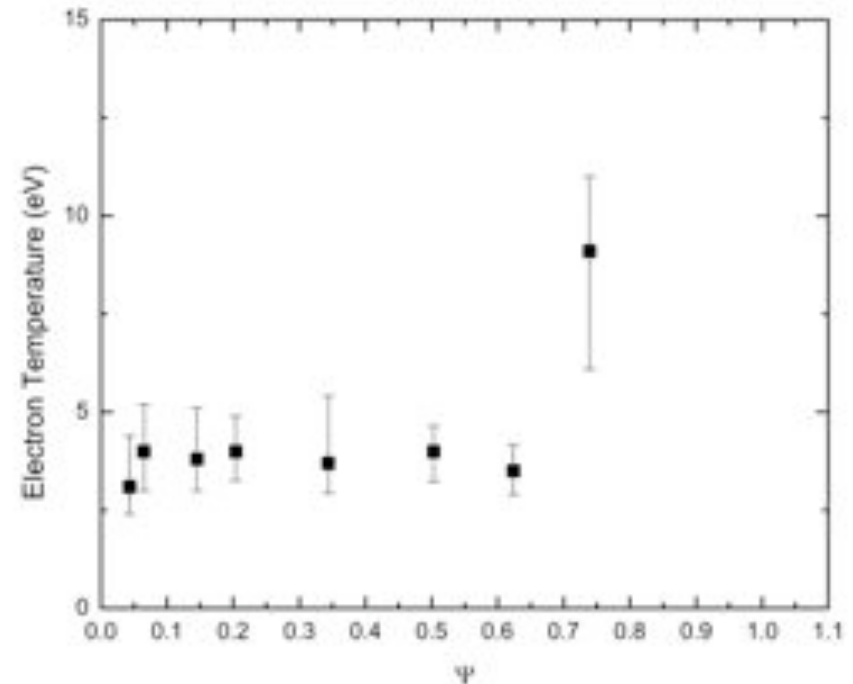
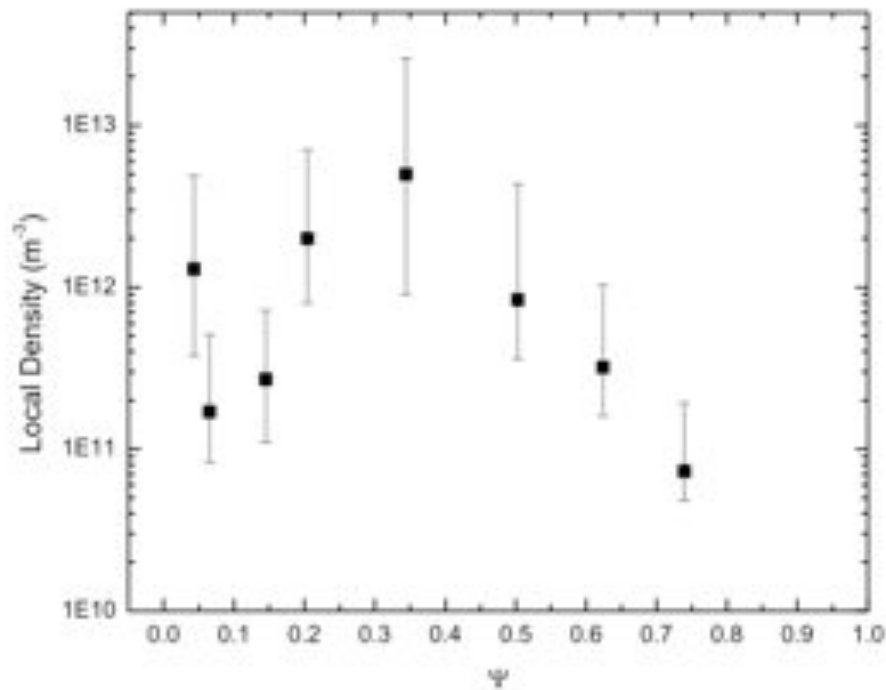
Filaments are halogen light bulbs (without the glass!)²

Each filament can serve multiple functions:

- Electron source
- Measure plasma potential
- Measure plasma temperature and density

² Inspired by M. Otte et al., IPP-Greifswald

Measured equilibrium density and temperature profiles



- Density measurements with probes in a pure electron plasma are challenging but possible (with large error bars)
- Interior temperature is flat at $\sim 4\text{eV}$
- Debye length $\approx 1.7\text{ cm} \ll$ minor radius of CNT - plasma criterion satisfied!

Equilibrium equation for a pure electron plasma

- Low density force balance

$$\vec{\nabla}P - en_e \nabla\Phi = \vec{j} \times \vec{B}$$

- **Perpendicular** force balance is trivial - gives perpendicular flow velocity (the B field is unperturbed by the plasma - low density)
- **Parallel** force balance yields Boltzmann distribution of the electron density on each magnetic surface
- This must be consistent with Poisson's equation:

$$\nabla^2\Phi = \frac{e}{\varepsilon_0} \underbrace{N(\psi) \exp\left(\frac{e\Phi}{T_e(\psi)}\right)}_{\text{density}}$$



Ψ is constant
on each surface
(magnetic flux
coordinate)

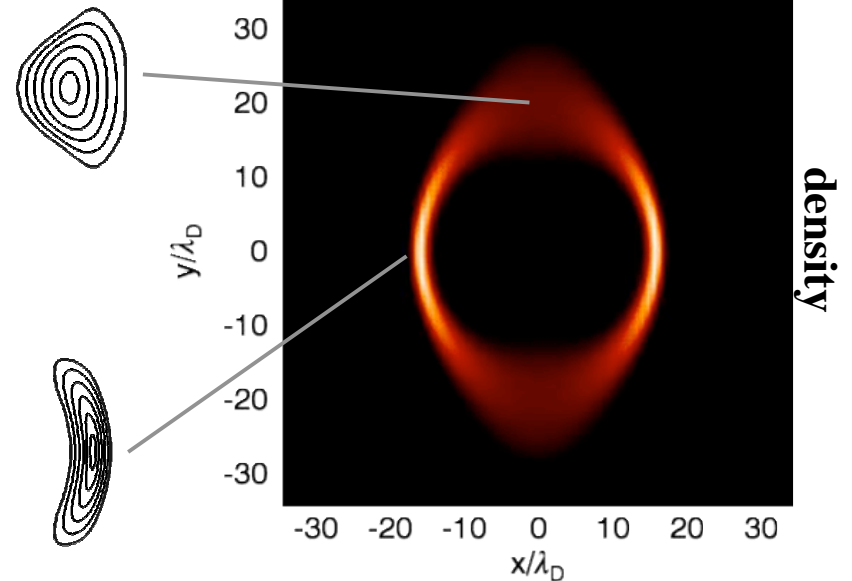
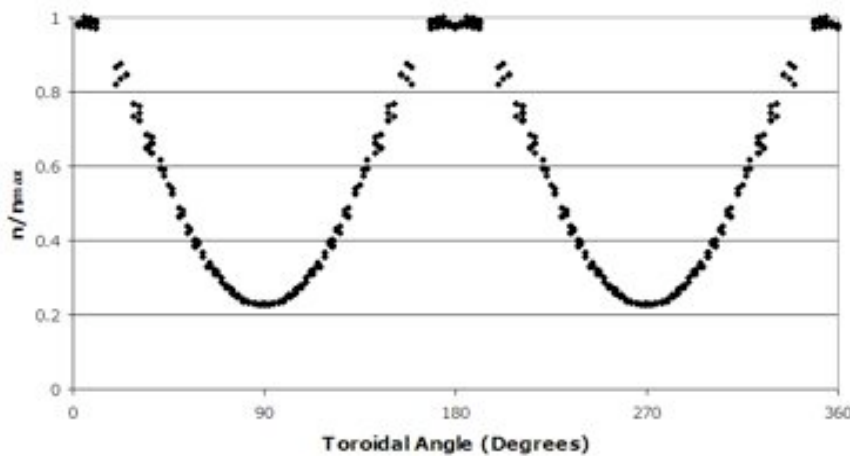
Equilibrium equation

$$\nabla^2 \Phi = \frac{e}{\epsilon_0} N(\psi) \exp \left(\frac{e\Phi}{T_e(\psi)} \right)$$

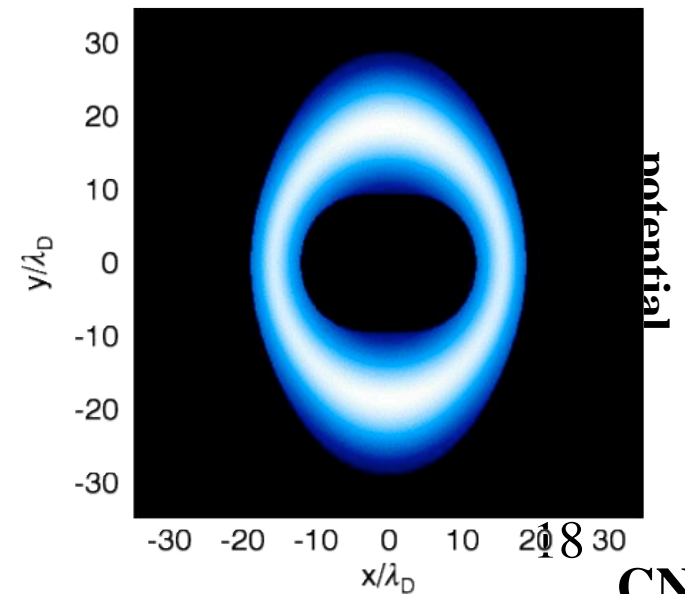
density

- Difficulty lies in non-linearity of problem
- Require two inputs to fully specify equilibrium:
 1. $N(\psi)$: provides freedom in density
 2. $T_e(\psi)$: temperature is assumed to be a flux function due to rapid parallel thermal transport (but may vary from surface to surface)
- A fully 3-D code has been developed to solve the equilibrium equation for arbitrary boundary conditions¹
- Measured density and temperature profiles allow a complete CNT equilibrium reconstruction

Toroidal density variation ($\lambda_D \sim 1.6$ cm)



- Large toroidal density variation due to changing cross-sectional area
 - Potential is roughly constant along field lines (for $\lambda_D \sim 1.6$ cm)
 - Density is higher at smaller cross-sections (like Penning-Malmberg / min B trap)
- Calculation predicts a factor of ~ 4 variation of density on the magnetic axis

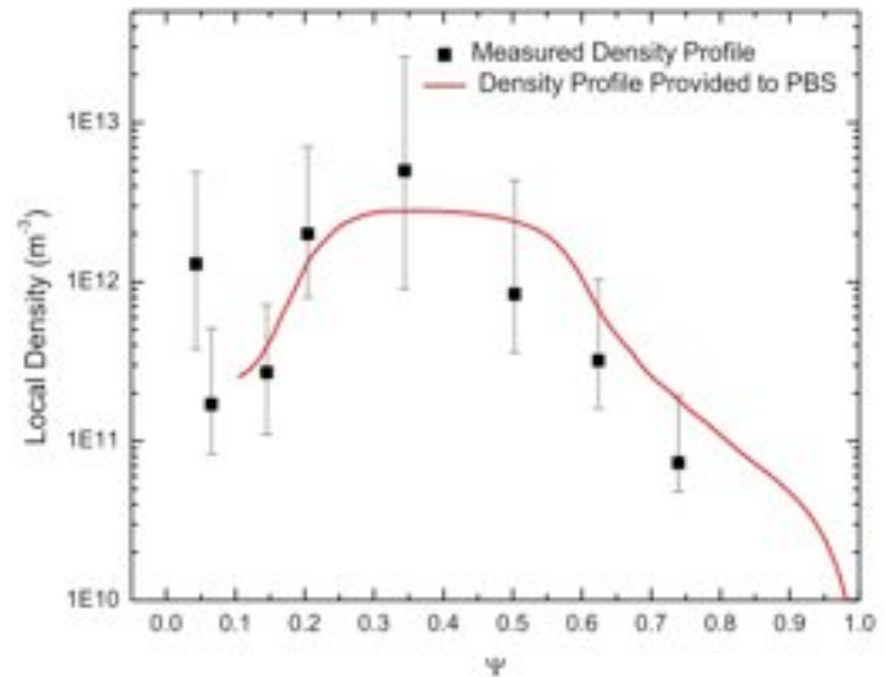
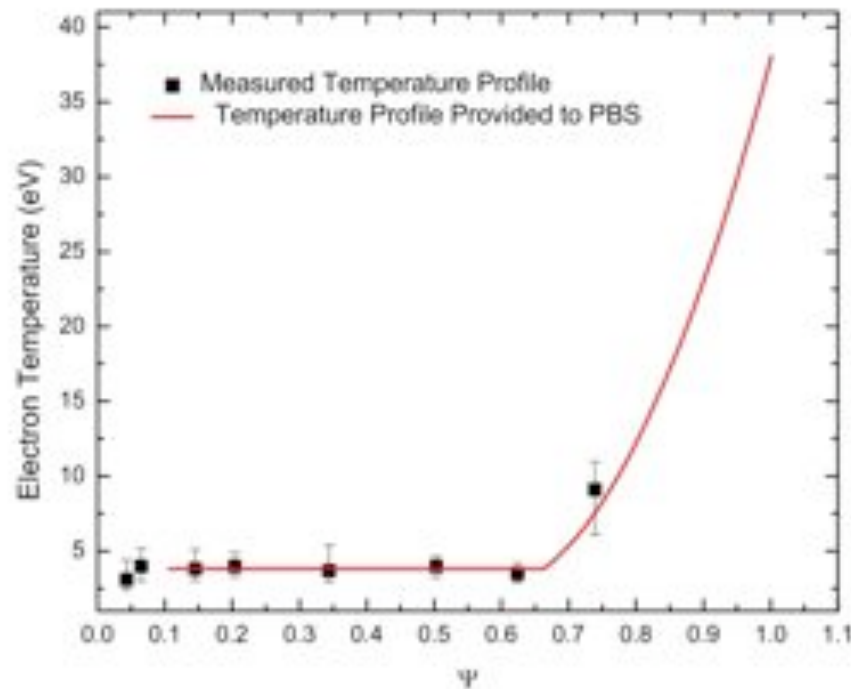


CNT

¹J. Fajans, Phys. Plasmas 10 p. 1209

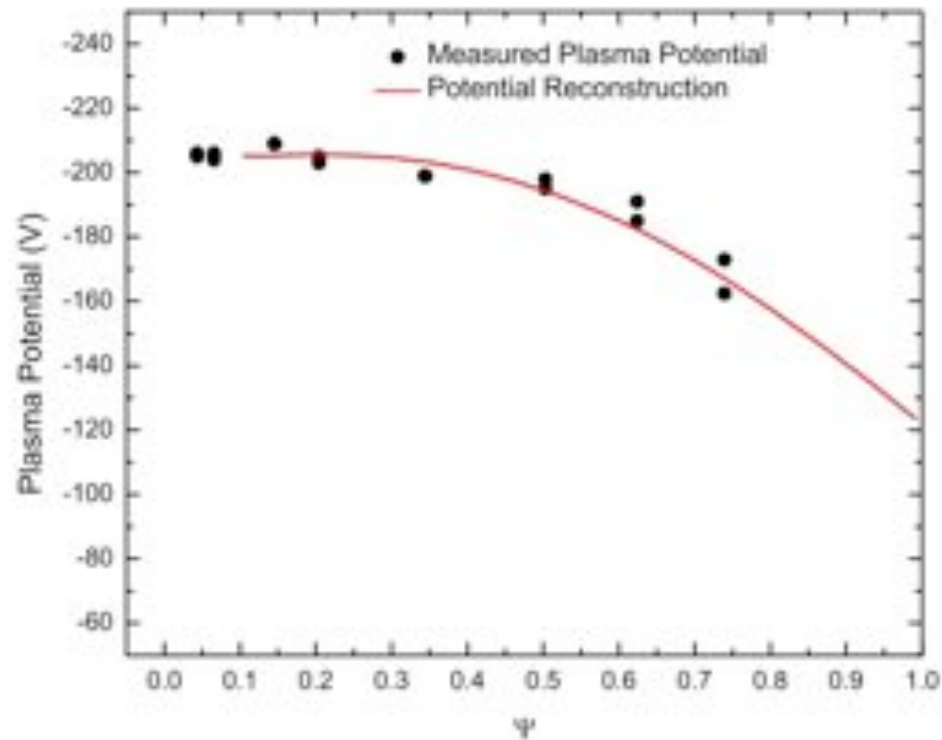
CNT

T_e and n_e profiles for equilibrium reconstruction



- Red graphs for density and temperature have been used to reconstruct CNT equilibrium, and calculate potential.
- Large temperature rise at edge is consistent with other measurements (not shown)

CNT equilibrium reconstruction



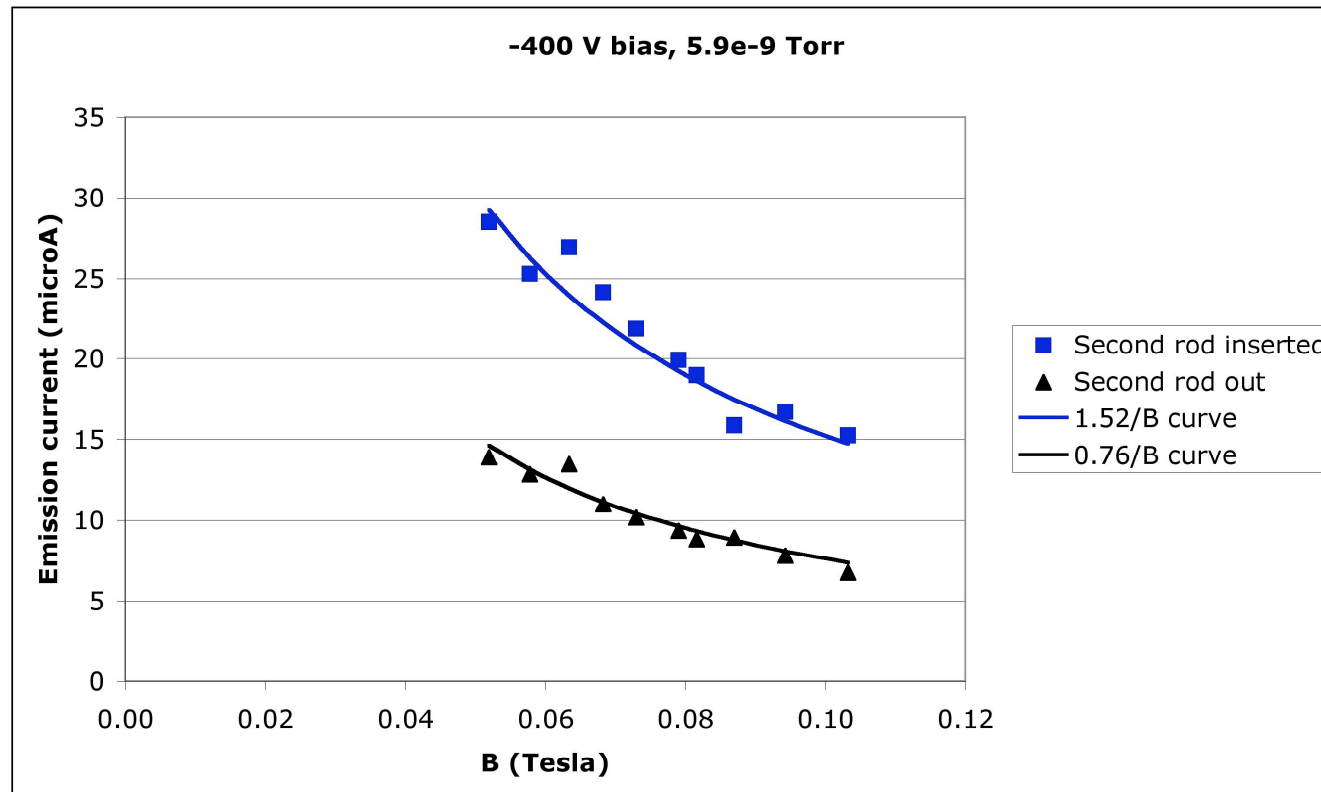
- Excellent agreement between calculated and measured potentials
- Equilibrium reconstruction also allows the electron inventory to be determined

Confinement time

- Knowing the electron source rate, and the total electron inventory, the confinement time can be computed. It's up to **~ 20 ms**
 - Parallel force balance (within a magnetic surface) establishes itself in ~ 10 μ s - electron fluid is in flux surface equilibrium
 - No large scale cross-surface $E \times B$ flows - these would lead to particle losses in ~ 10 - 100 μ s
 - No significant direct "bad orbit" losses - direct ∇B drift out of the device would occur in ~ 300 μ s
- **Conclusion: We have a macroscopically stable equilibrium**
- What limits the confinement? Why is it not many seconds?



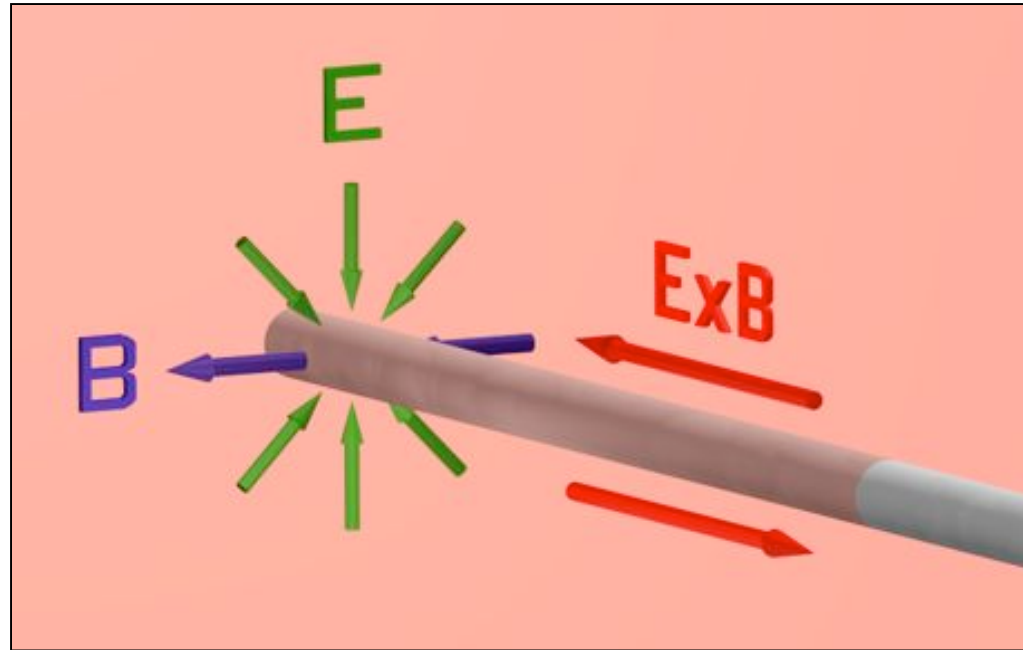
Insulated rods limit confinement



- Two rods gives twice as much transport as one rod (for low neutral pressure and relatively large bias voltage)
- The rods are insulating, so they are not steady state sinks for electrons
- They are large electrostatic perturbations - drives ExB transport



Insulated rods limit confinement

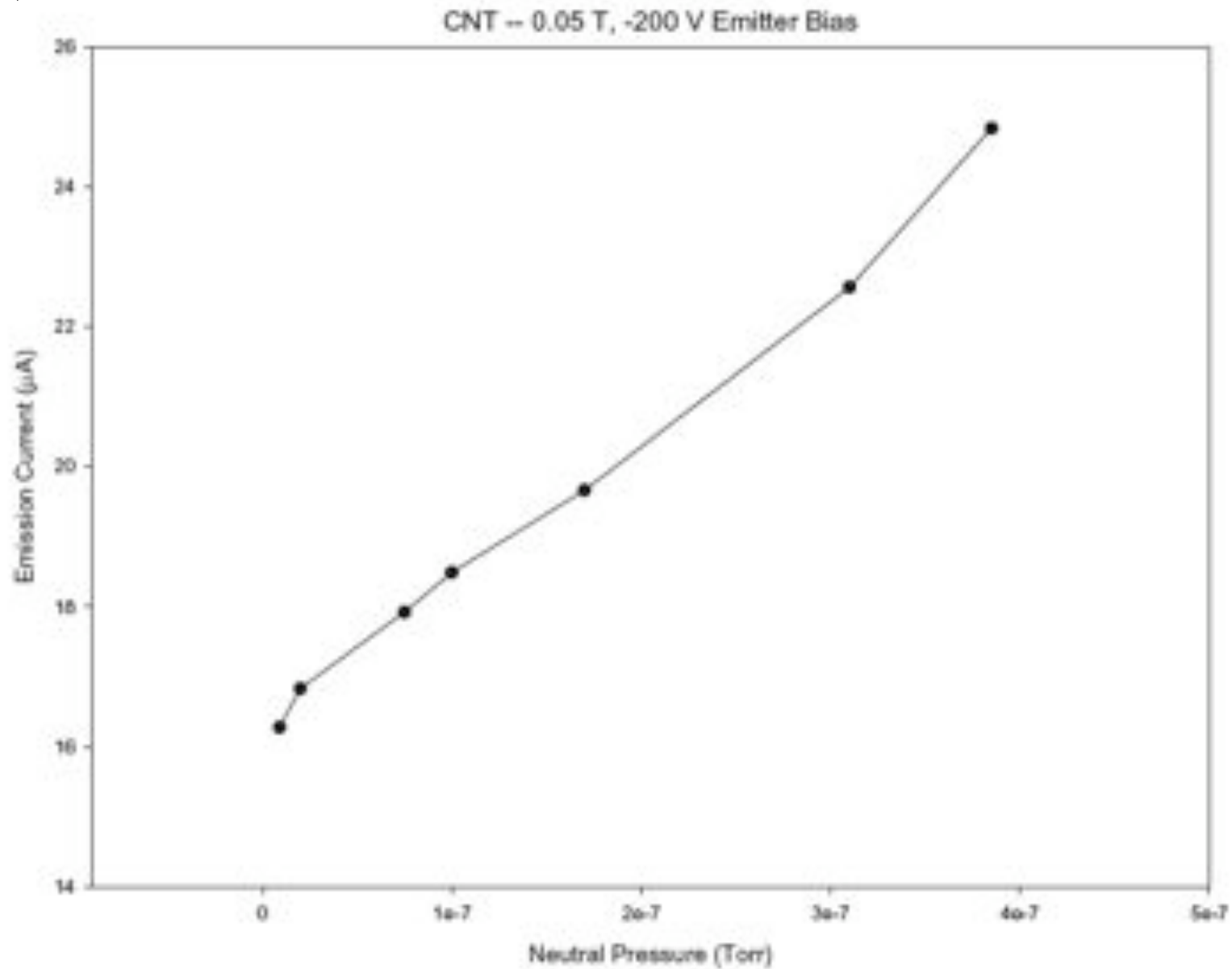


Insulated rods charge up negative relative to plasma

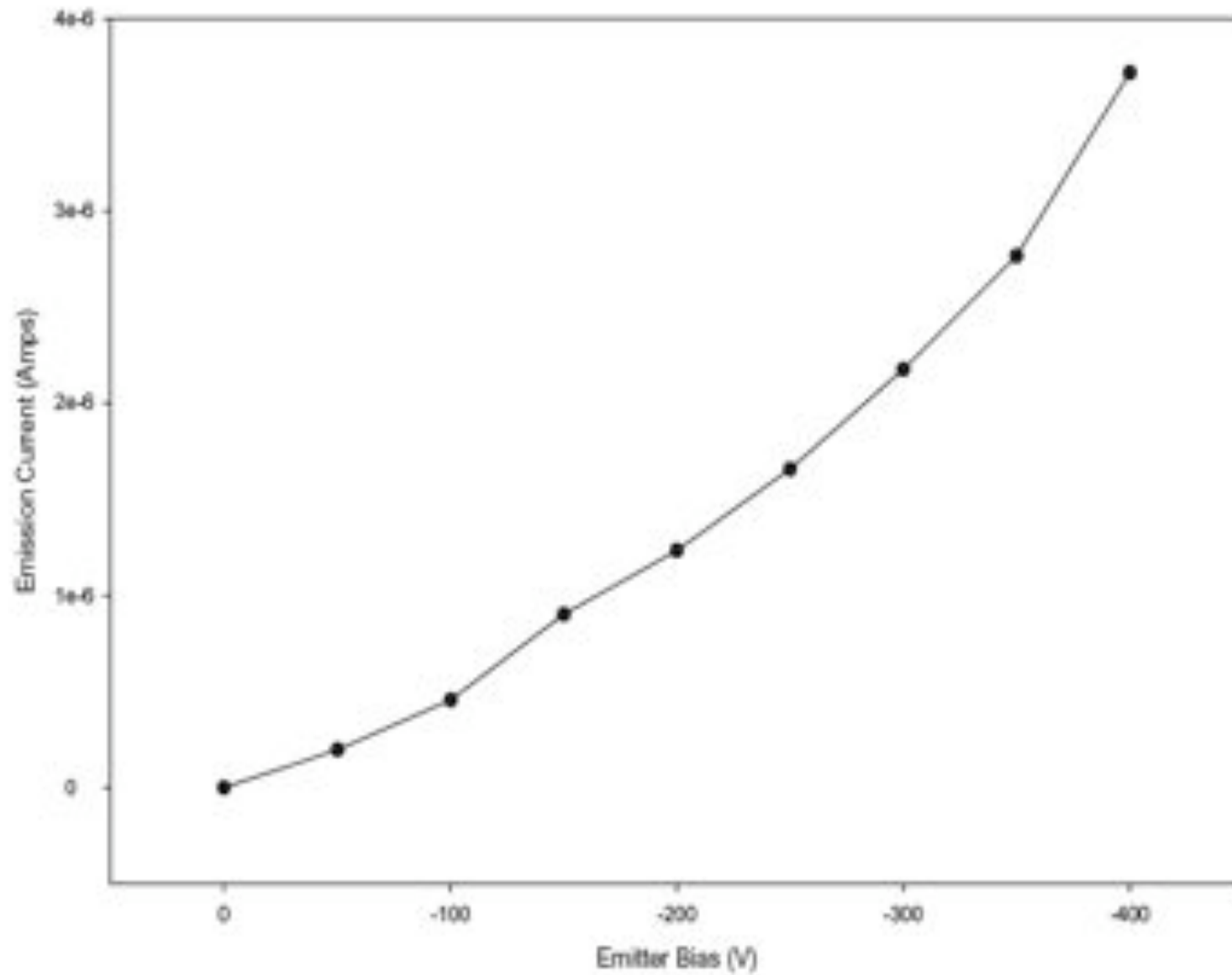
Resulting $E \times B$ drift pattern convects particles along the rod all the way to the open field lines

Observed order of magnitude and $1/B$ scaling of transport rate is consistent with this effect

Neutrals also degrade confinement



Bias voltage affects confinement



Retractable emitter - eliminate rod driven transport

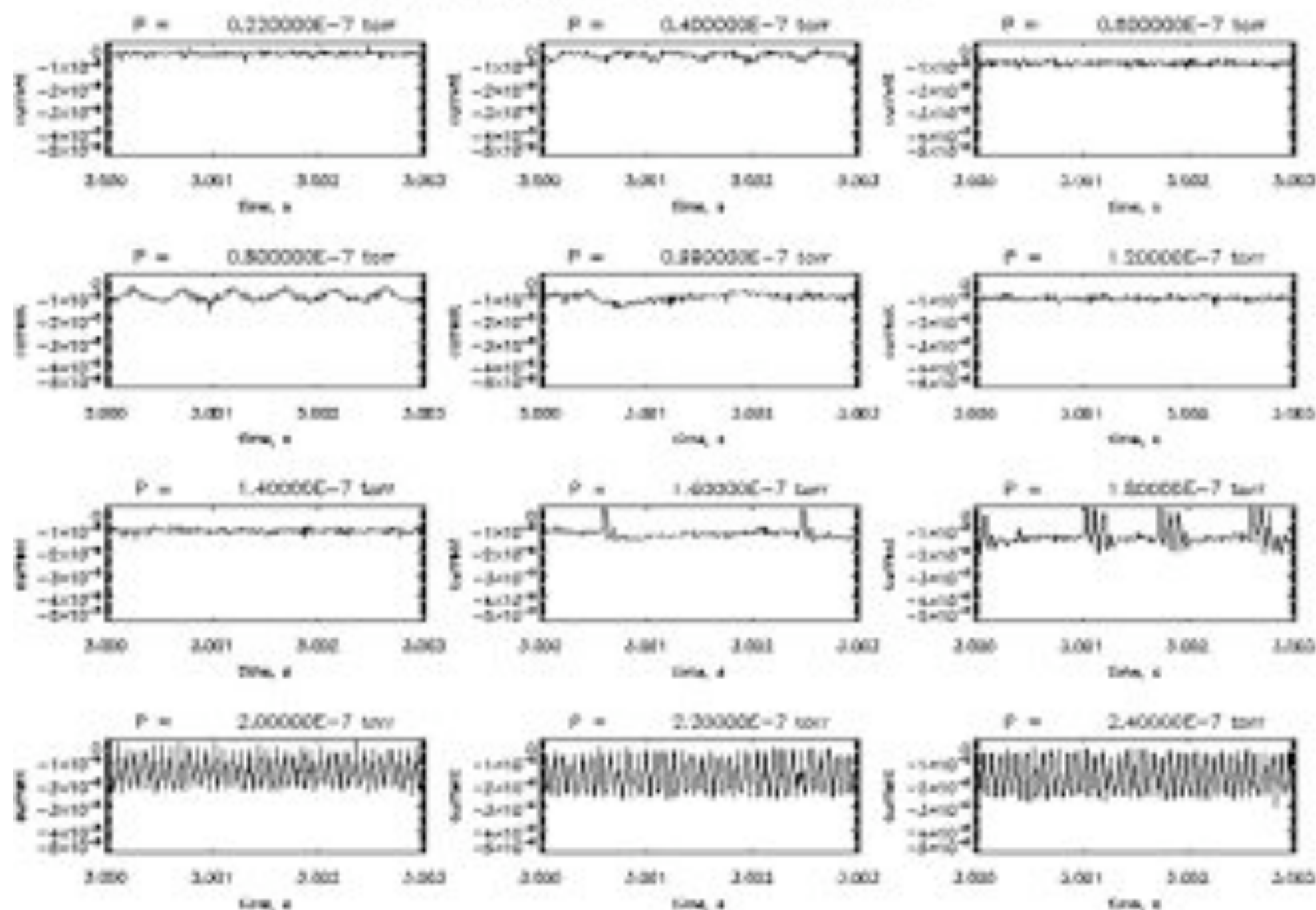


Ions

- At neutral pressures of $5 \cdot 10^{-9}$ - $1 \cdot 10^{-8}$ Torr the ion content is $<1\%$
 - Essentially a pure electron plasma
 - Measured with large probe inserted specifically for this purpose
- If there were no sink then ions would continually accumulate - they don't
- Insulated (negatively charged) rods act as sinks for ions - allow a steady state ion density to establish itself
- At much higher neutral pressures ($\sim 2 \cdot 10^{-7}$ Torr) , CNT has 4-10% ions. At relatively low B-field strengths, this triggers an instability, similar to what has been seen in the CHS stellarator.¹



Emission Current vs. Neutral Pressure, -200 V, .02 Tesla



Characteristic frequencies

- Observed oscillations are in the 10-50 kHz range
- For $\phi_{\text{bias}}=200$ V, $B=0.02$ Tesla plasmas, assuming singly ionized nitrogen atoms:
 - ExB drift velocity ~ 20 kHz
 - Ion Larmor frequency ~ 20 kHz
 - Ion radial oscillation frequency (ignoring B) ~ 30 kHz
- More work planned to understand these oscillations:
 - Use different ion species (change m_i)
 - Measure spatial structure (new probes and electronics)
 - Compare with theories and observations of ion-resonant instabilities in other confinement devices^{1,2,3}

¹R. H. Levy et al, Phys. Fluids 12. P 2616

²A.J. Peurrung et al., PRL 70 p. 295

³M. R. Stoneking et al. Phys. Plasmas 9 p. 766



Near term plans

- Measure confinement time and plasma behavior in the absence of rods
 - Will ions continually accumulate?
 - Will ion instability destroy plasma?
 - What will be the confinement time and confinement scalings?
- More spatially resolved equilibrium measurements
 - Full 2-D cross sectional measurements of n_e , T_e , plasma potential
 - Measurements at two principal toroidal cross sections



Near term plans

- Investigate physics of ion driven instability
 - Spatially resolved measurements
 - Determine mode structure (toroidal and poloidal)
 - Determine radial profile of instability amplitude
- Equilibrium:
 - Radial profile control
 - inject on multiple surfaces simultaneously
 - tailor potential profile - can we affect the temperature? Can we change the density profiles? Can we change confinement?



Not so near term plans

- Explore the two other tilt angles in CNT
 - 39° tilt angle: No shear, moderate iota
 - 44° tilt angle: Large iota, extreme shaping
- Electron-positron plasma research
 - Non-neutral stellarator may be ideal device for the creation and studies of electron-positron plasmas
 - Proposed injection scheme:
 - Create a well confined electron plasma
 - Open up stochastic edge region
 - Inject positrons on a stochastic (but long) field line



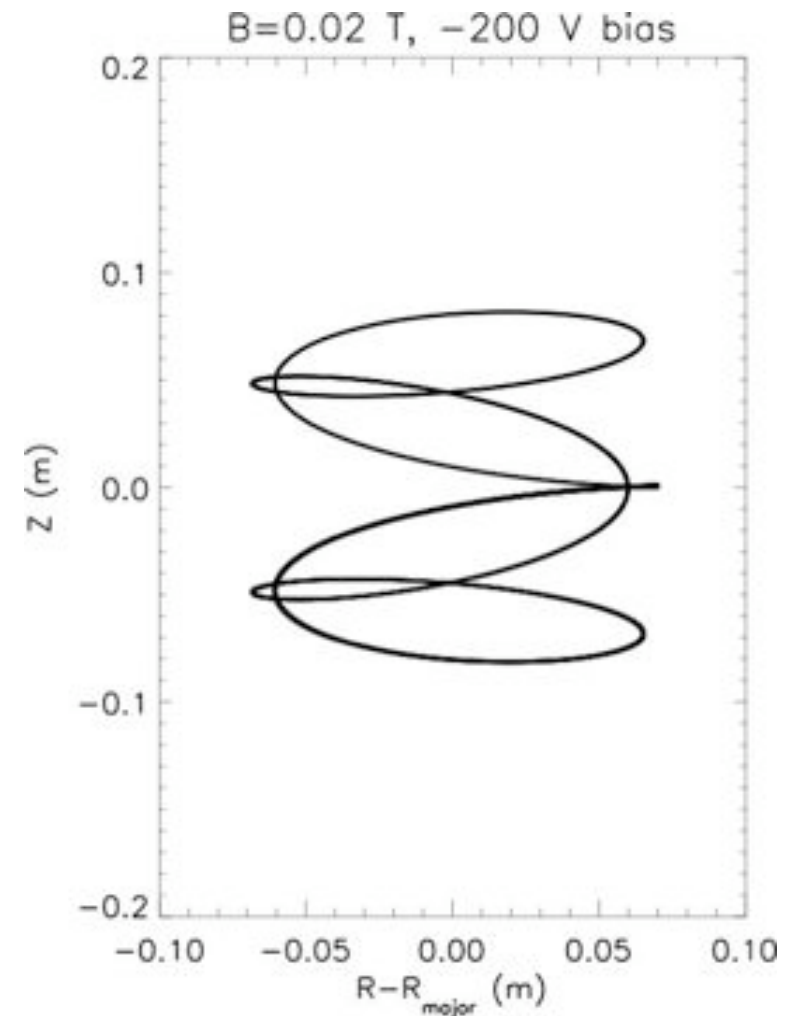
Theory opportunities: Overview

- What is the physics of the ion instability in CNT?
- Why is confinement so sensitive to neutral pressure in CNT?
- Why does the neutral driven transport scale as $B^{-1.5}$?
- Can we understand the profiles of density and temperature in CNT?
- Can a CNT-like configuration be fusion-relevant?

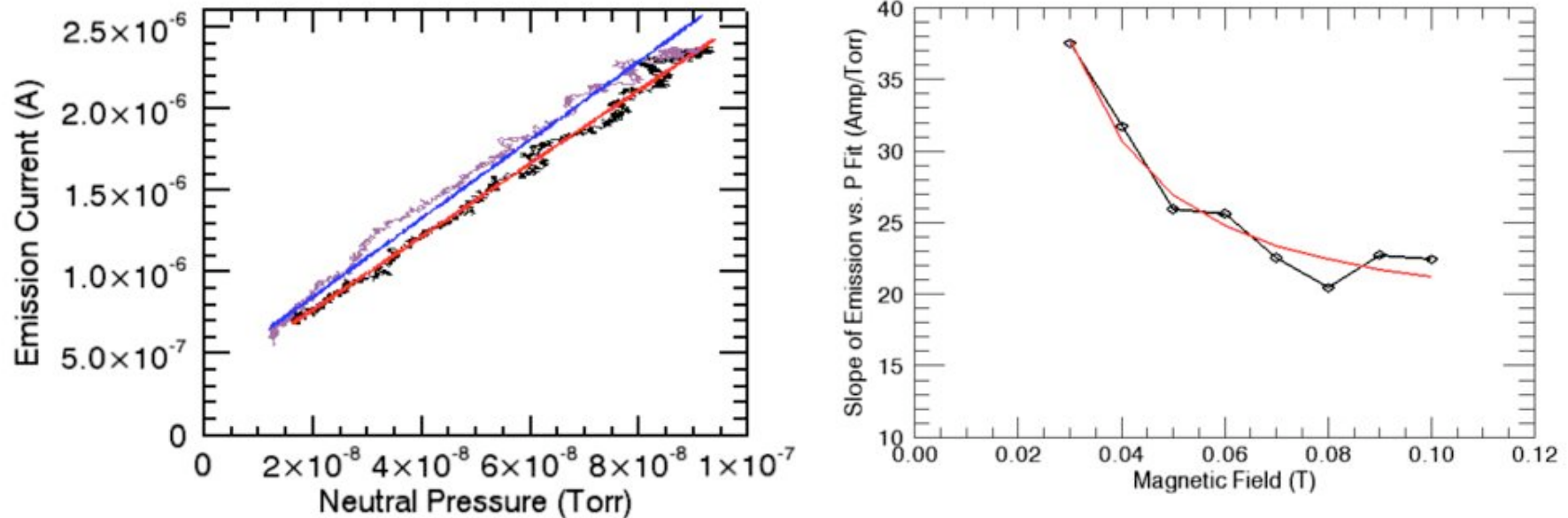


Theory opportunities: Ion instability

- Develop a theory for the ion instability in a magnetic surface device
 - Ion instability is well understood in Penning traps
 - Electron dynamics are fundamentally different in a stellarator - electron perturbation should damp out, except on rational surfaces
 - Ions need to be treated kinetically - particle orbits are on the order of the radial dimension



Transport due to neutrals



Confinement is determined by loss due to all sources (rods, classical, neo-classical...). The rods are a major source but other sources are significant:

- Emission scales linearly with p_n .
 - loss rates are much larger than neoclassical diffusion
- Slopes of Emission vs. Neutral pressure scale as $\sim B^{-1.5}$
 - Same scaling has been observed in Stoneking's pure toroidal field experiment
- Are there still direct loss orbits in CNT despite the large electric field?

Theory opportunities: Profile effects

- Understand measured n_e and T_e profiles
 - Why do we see a peak off axis in density even though we are injecting at the axis?
 - Implies a non-diffusive transport mechanism
 - Implies good local confinement in region of bump
 - Why is the electron temperature flat throughout most of the plasma? (electron - electron collision time is $>$ confinement time)
 - One would expect “Joule heating” - electrons should pick up kinetic energy from the electrostatic potential as they diffuse out
 - Why does the temperature rise steeply near the edge?
 - Confinement is poor at the edge (we see low density) - so perhaps particles can Joule heat and not interact with the rest of the plasma?



Theory opportunities: Fusion relevance of CNT?

- CNT has several advantages as a fusion device, all of which have been verified experimentally:
 - Ultralow aspect ratio (compact device)
 - Extremely simple coils
 - Error field resilience
- But it is not a good candidate for fusion:
 - Far from any quasisymmetry - predicted very high neoclassical losses (H. Mynick, N. Pomphrey, PPPL, Seiwald, U. Graz)
- Can we optimize a CNT-like configuration for fusion, retaining the advantages of CNT?
 - Just started neoclassical optimization effort with B. Seiwald, Tech. Univ. of Graz, Austria



Summary

- CNT can explore unique non-neutral plasma physics
- Stable equilibria exist for pure electron plasmas in a stellarator
- Small Debye length has been achieved in CNT
- Confinement is decent - limited by presence of rods and neutrals
- Instabilities can be triggered by poor vacuum, large bias potential, or low magnetic field
- Plenty of phenomena to study!

More info: Contact me at tsp22@columbia.edu or look on our website: www.ap.columbia.edu/CNT

